

Optical Properties of Solids: Lecture 12

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Optical Properties of Solids: Lecture 12

Quantum structures: wells, wires, dots

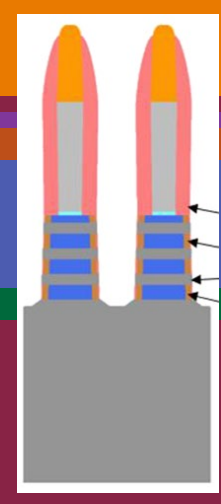
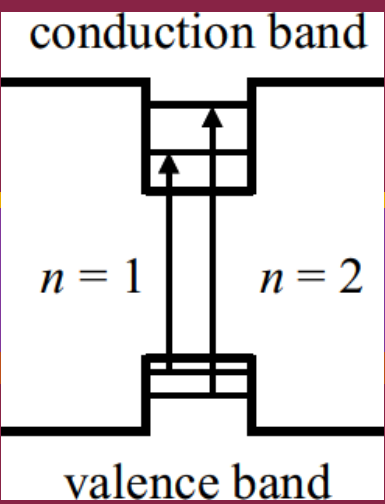
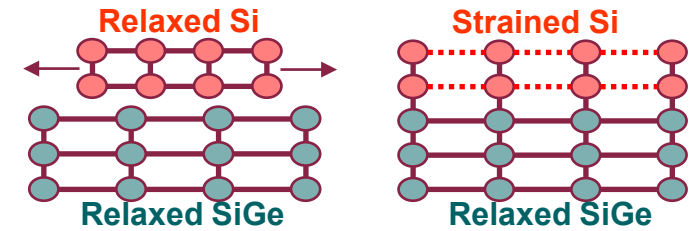
Absorption and emission of quantum wells

Intersubband absorption

Quantum cascade lasers

Defects (deep, shallow)

Stress and strain



References: Band Structure and Optical Properties

Solid-State Theory and Semiconductor Band Structures:

- **Mark Fox, Optical Properties of Solids (6,8,9)**
- Ashcroft and Mermin, Solid-State Physics
- **Yu and Cardona, Fundamentals of Semiconductors**
- Dresselhaus/Dresselhaus/Cronin/Gomes, Solid State Properties
- Cohen and Chelikowsky, Electronic Structure and Optical Properties
- Klingshirn, Semiconductor Optics
- Grundmann, Physics of Semiconductors
- Ioffe Institute web site: NSM Archive
<http://www.ioffe.ru/SVA/NSM/Semicond/index.html>

Outline

Quantum confinement and Heisenberg uncertainty principle

Growth of quantum structures

Carbon nanostructures, two-dimensional materials

Electronic states, quantum well absorption and emission

Intersubband transitions

Metamaterials and metasurfaces

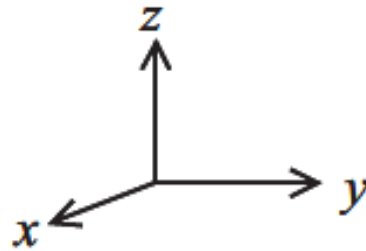
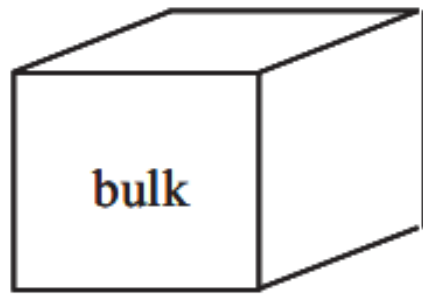
Defects

Transition metal and rare earth impurities in insulators

Shallow defects in semiconductors

Stress and strain, deformation potentials

Quantum structures and density of states



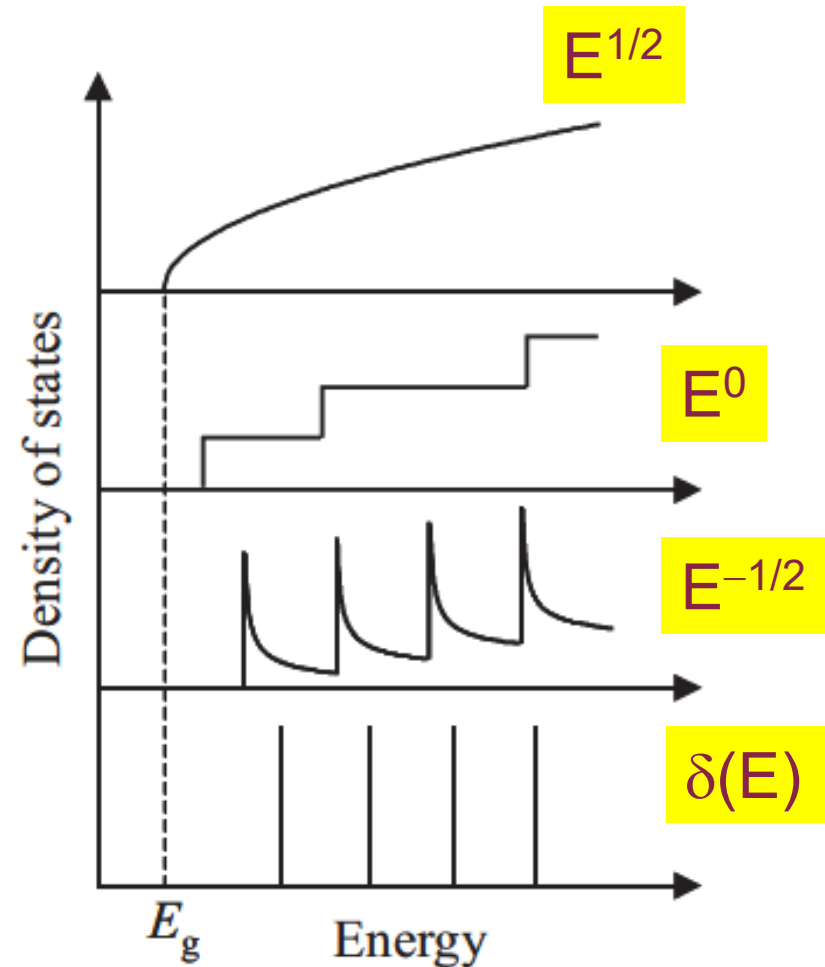
quantum well



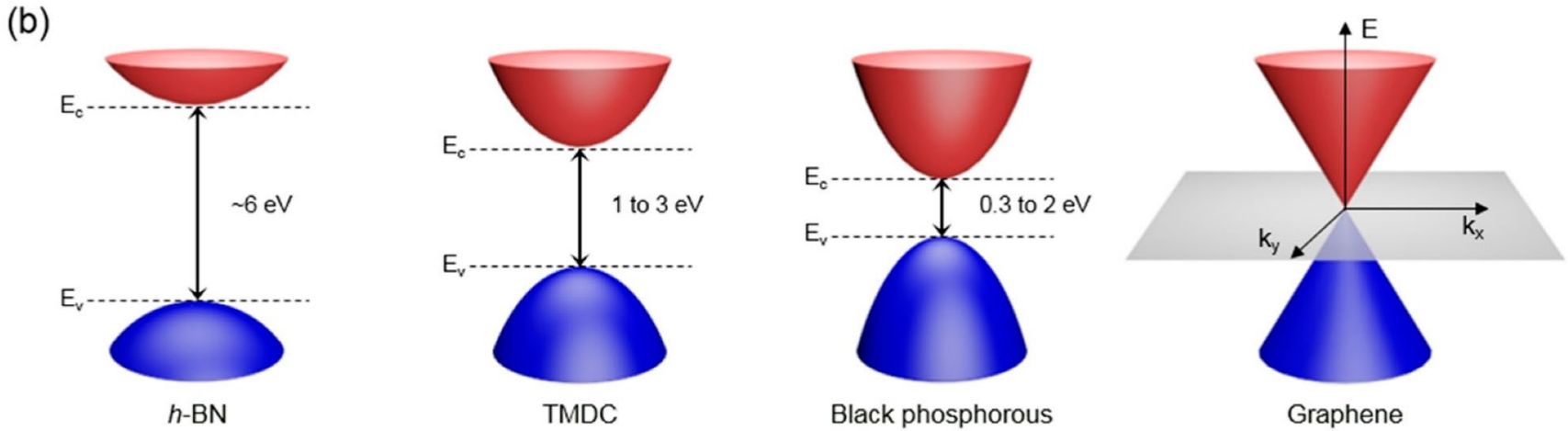
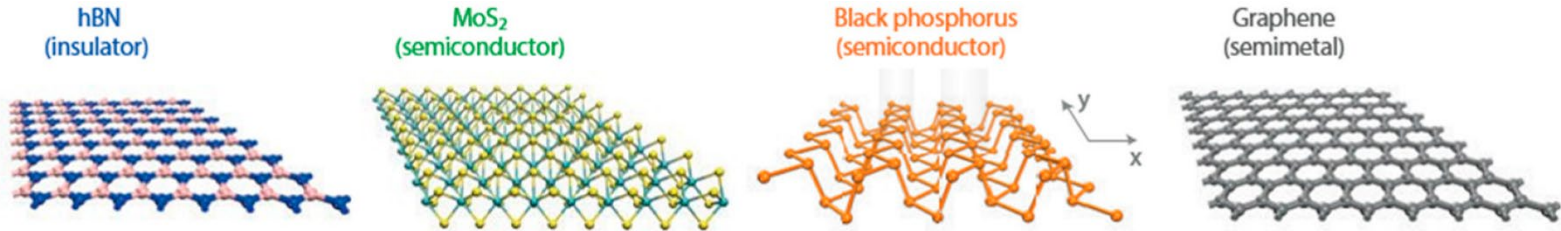
quantum wire



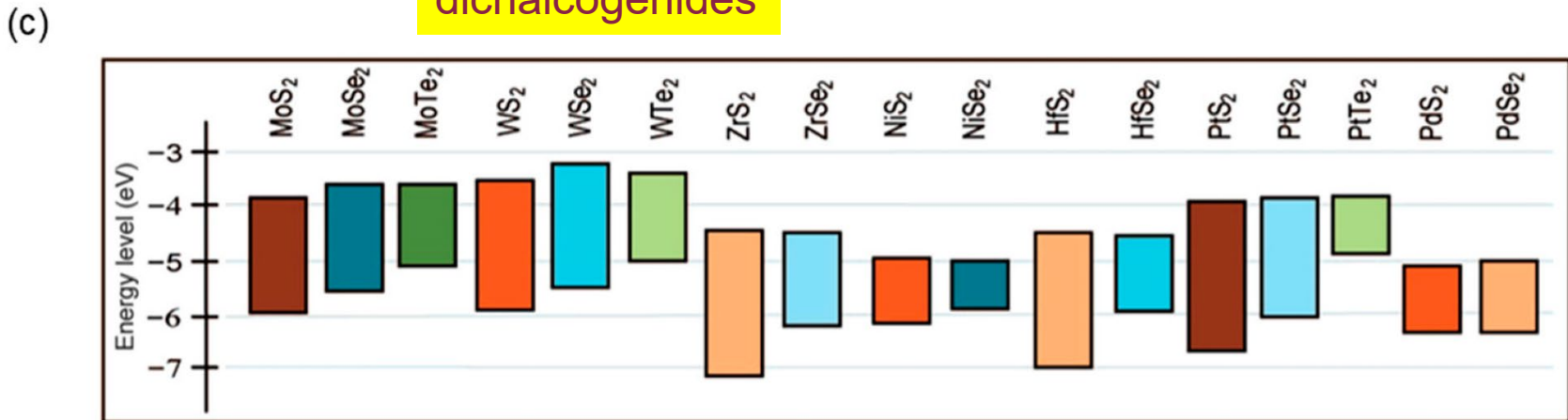
quantum dot



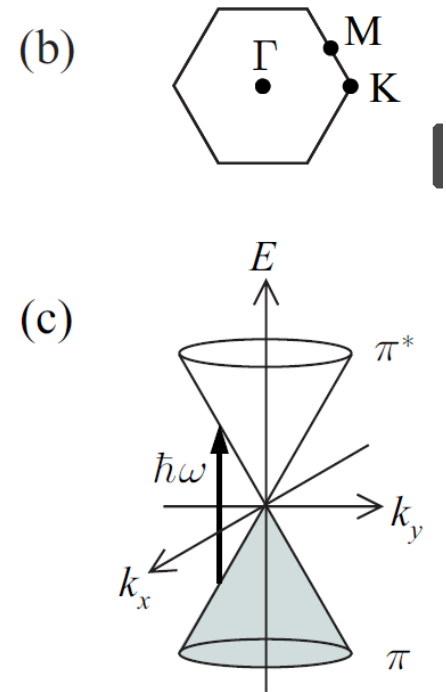
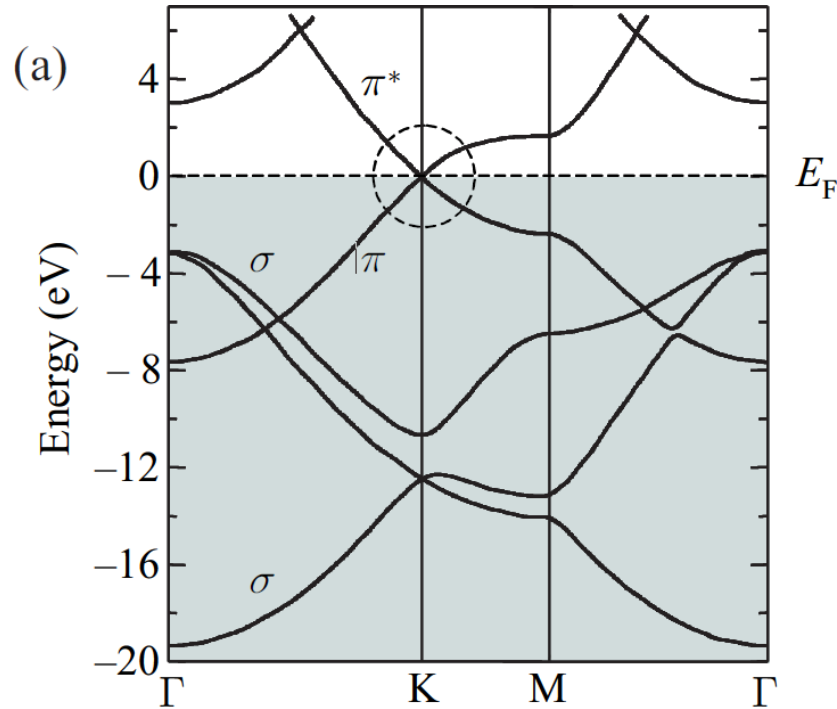
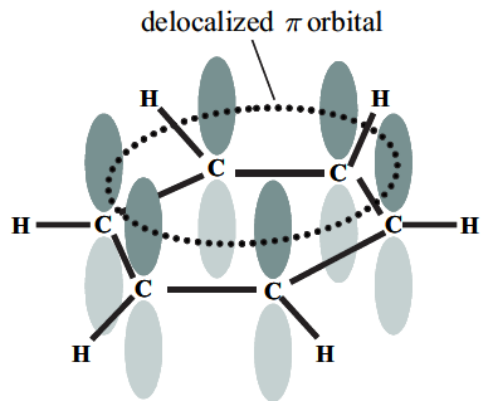
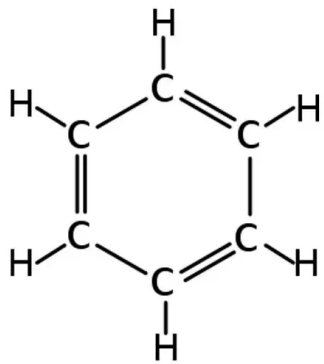
Two-dimensional semiconductors



dichalcogenides

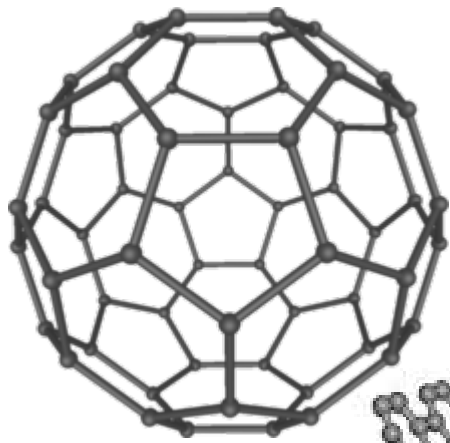


Dirac point in graphene

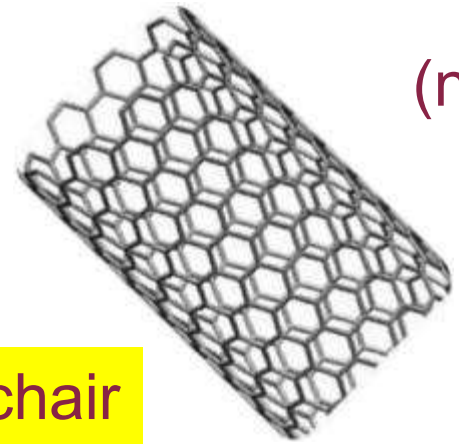
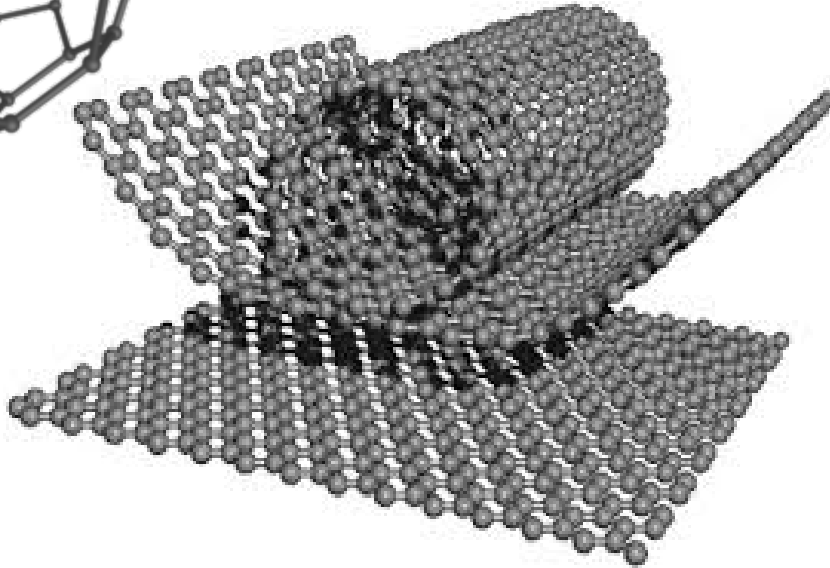


Graphite: Strong bond in the graphene plane, weak bond between planes. Remove single graphene plane.

Carbon nanotubes and fullerenes

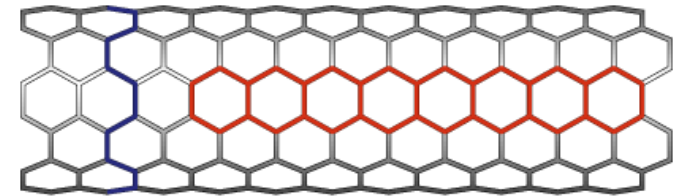


C-60

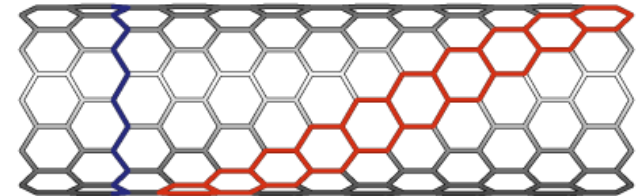


(n,m)

armchair



Armchair



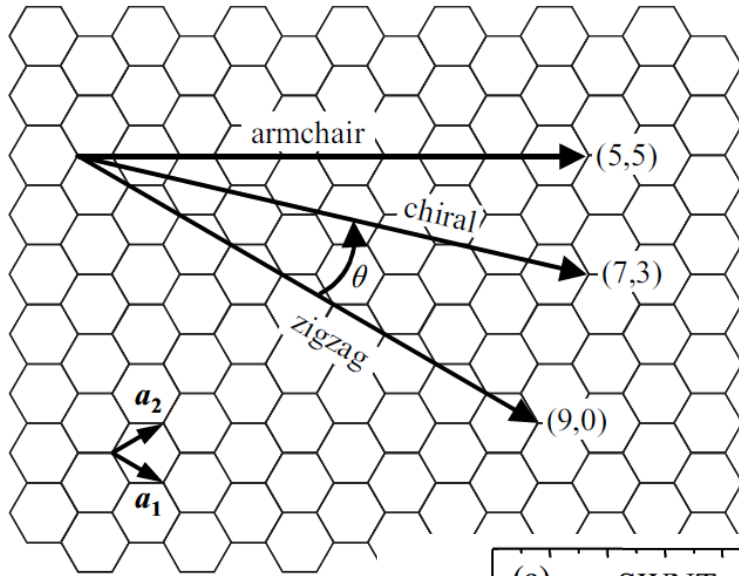
Zig-zag

Zig-zag

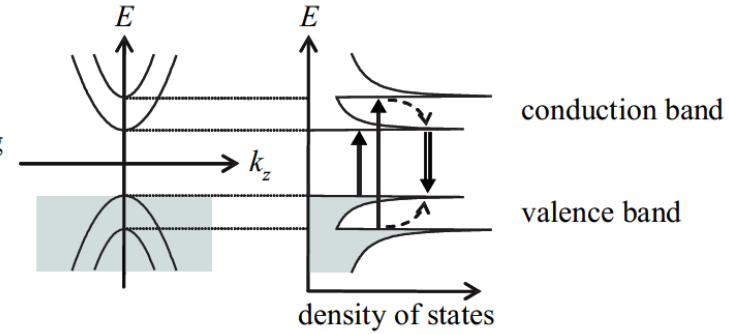
Graphene sheets are rolled up to form tubes.
Many Raman studies, doping with defects.

(n,m)

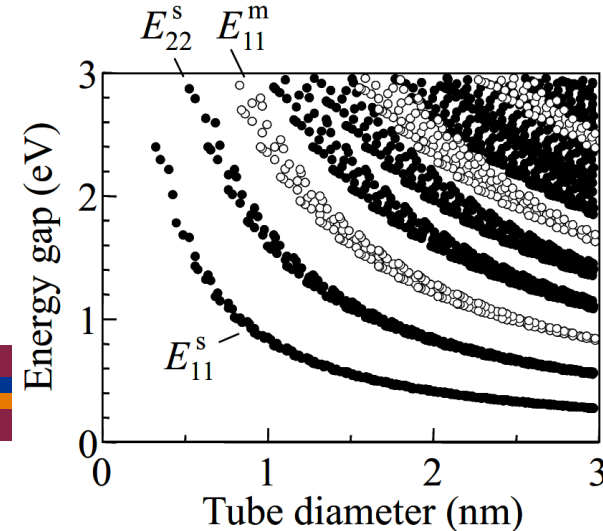
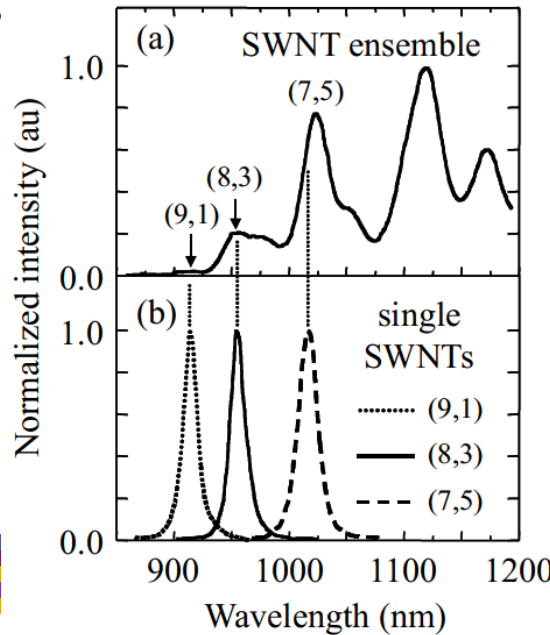
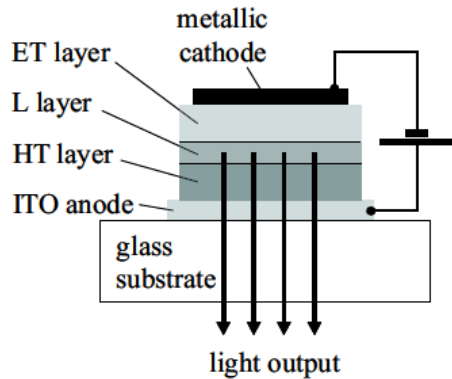
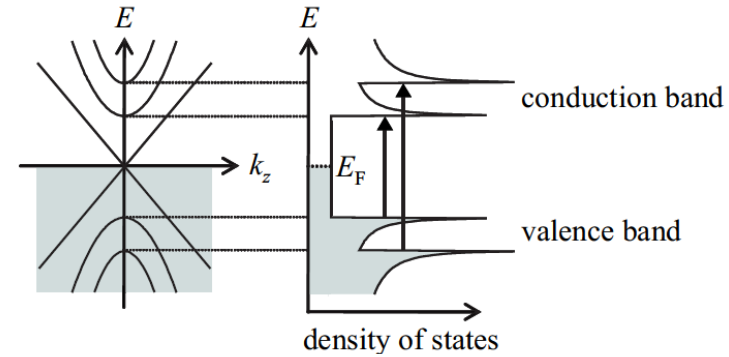
Carbon nanotubes: electrical properties



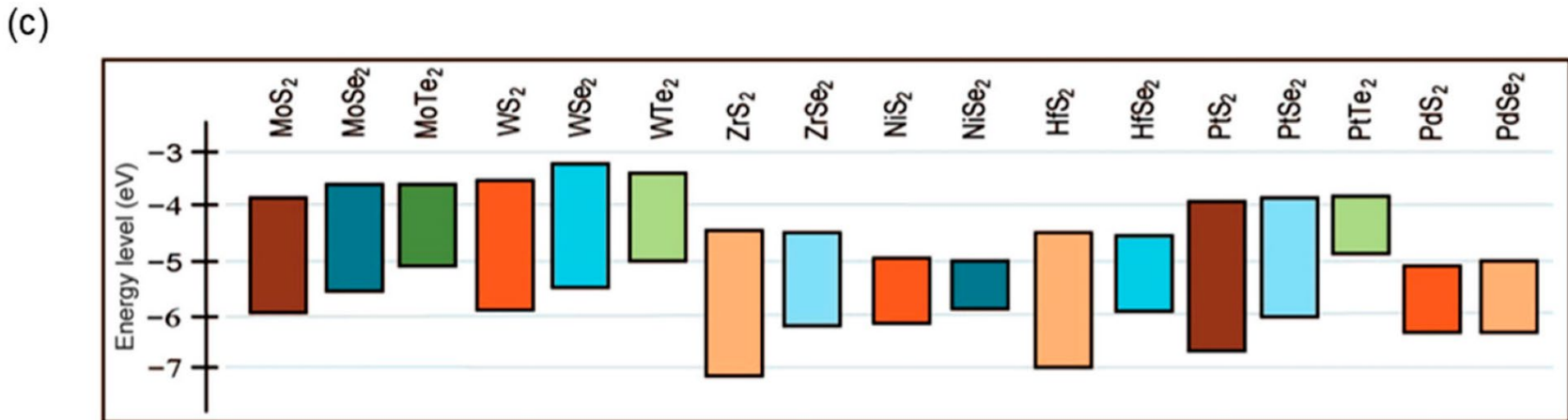
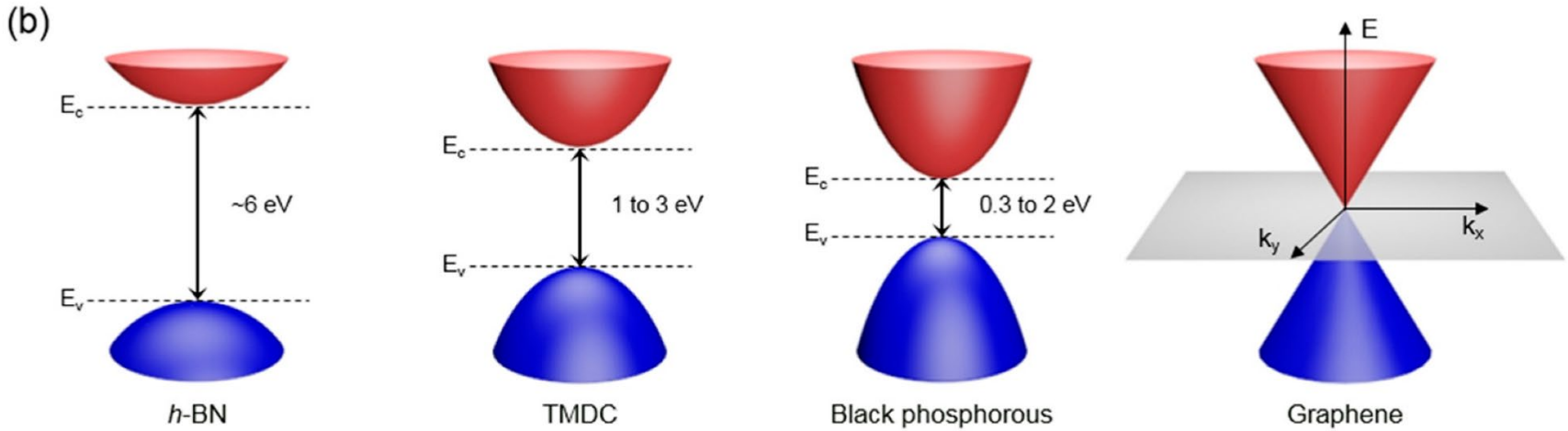
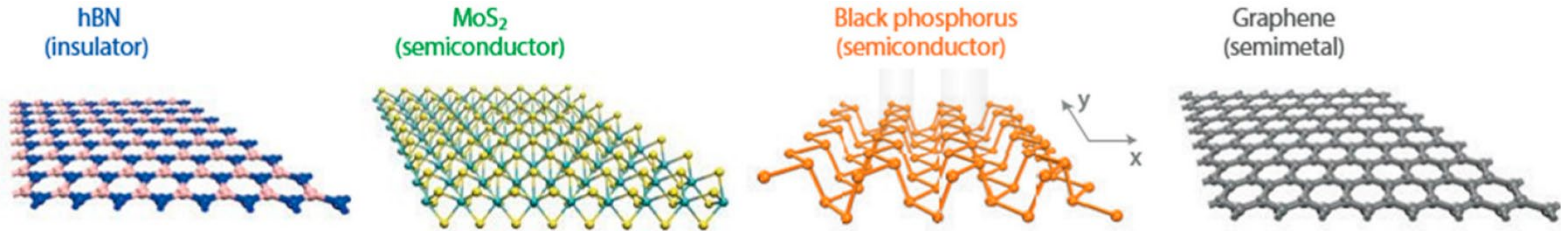
(a) semiconducting nanotube



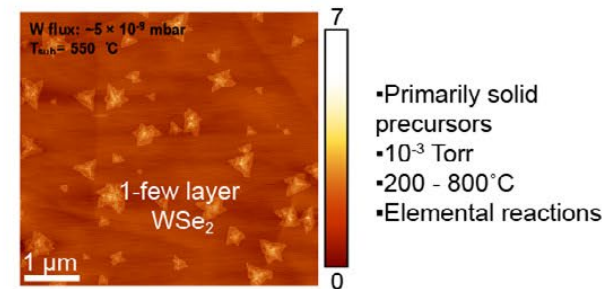
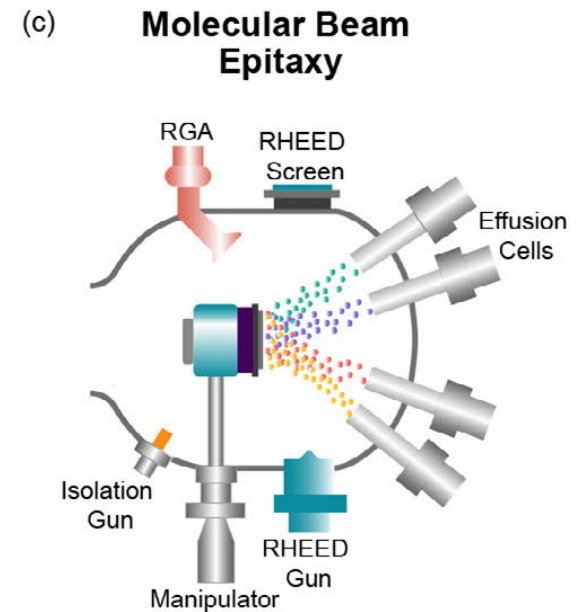
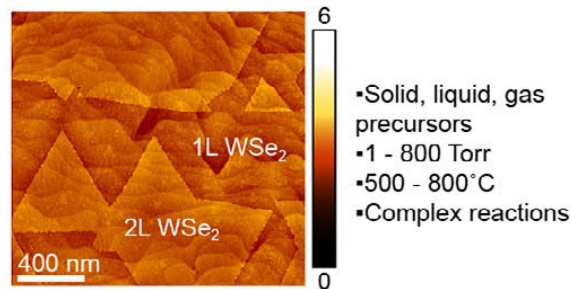
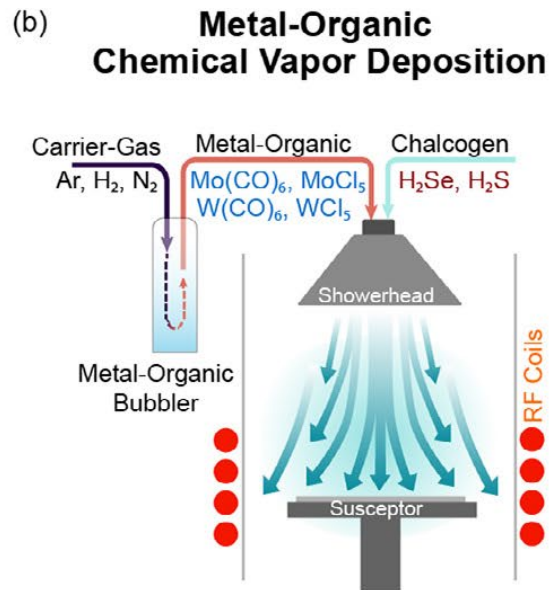
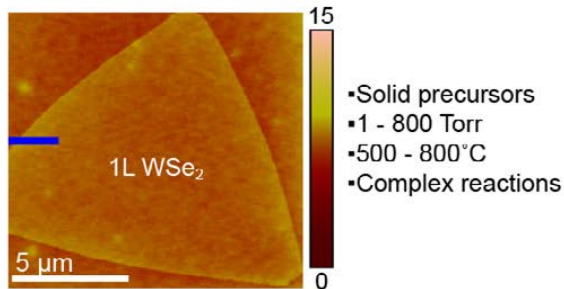
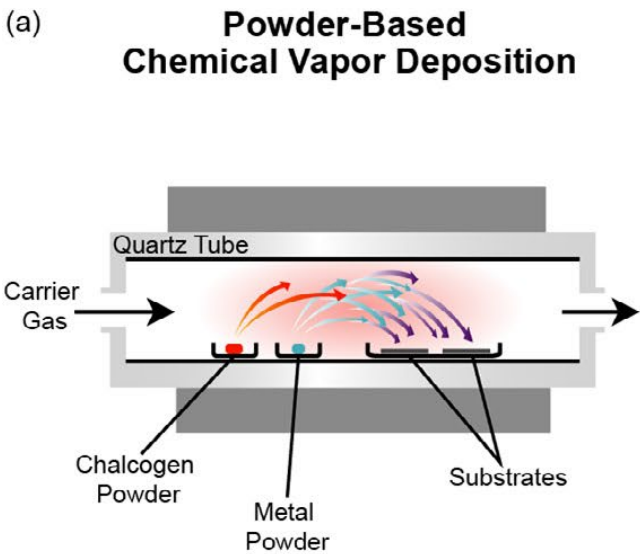
(b) metallic nanotube



Two-dimensional semiconductors

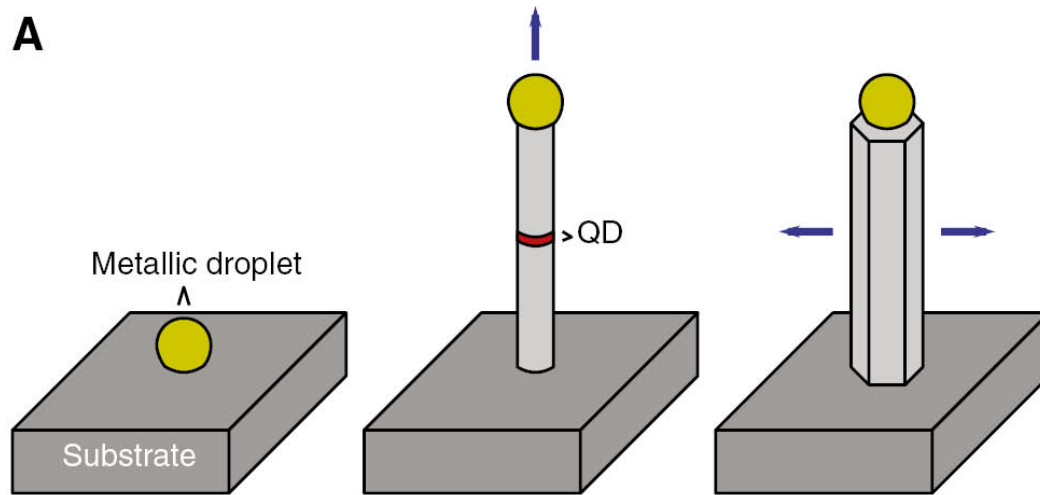


Two-dimensional semiconductors: WSe_2 , MoSe_2



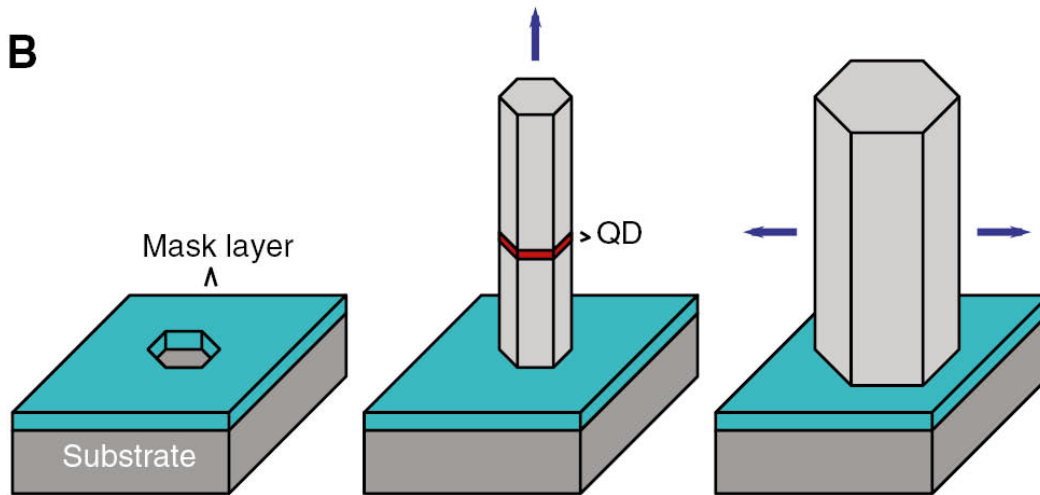
Growth of vertical quantum wires

A



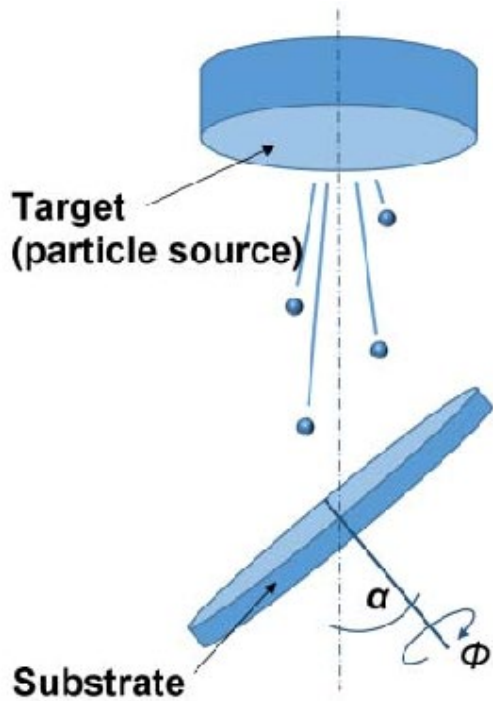
Vapor-liquid-solid

B

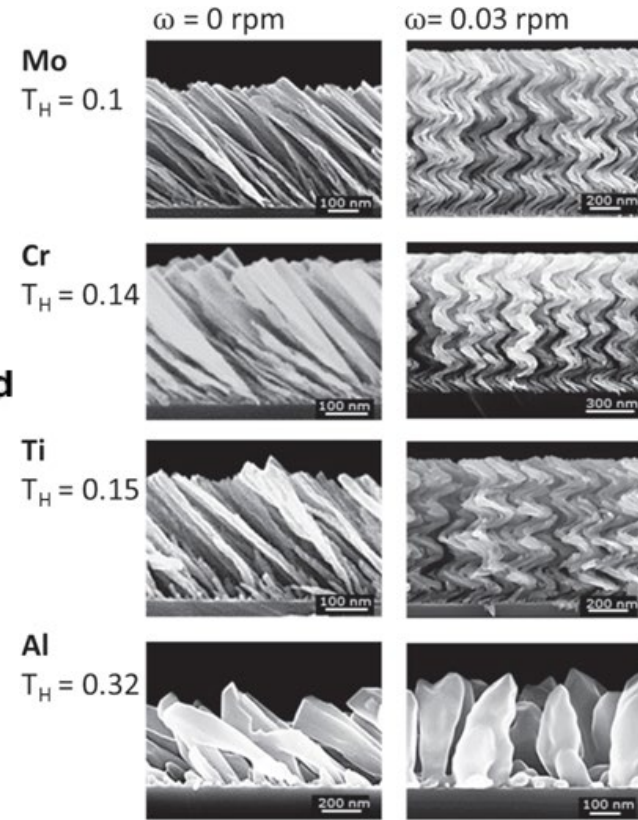
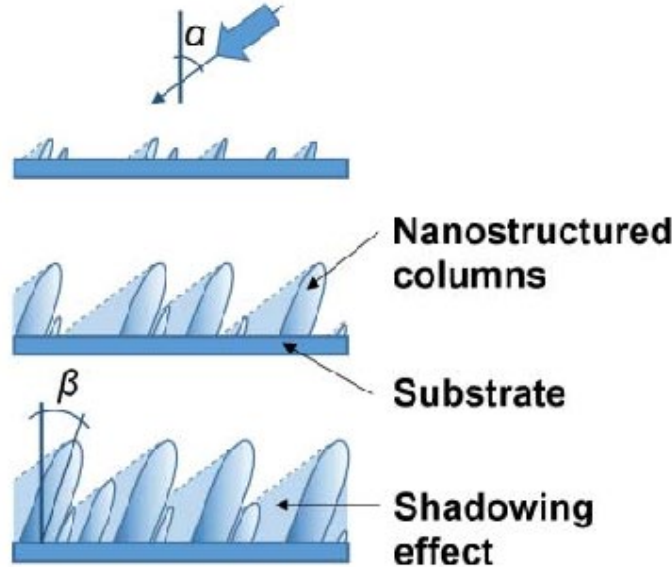


Selective area epitaxy

Glancing angle deposition (GLAD) of quantum wires



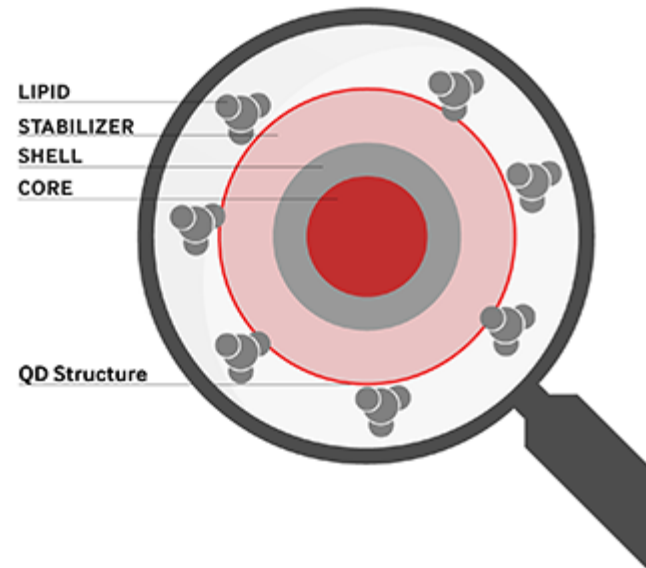
Flow of incident particles



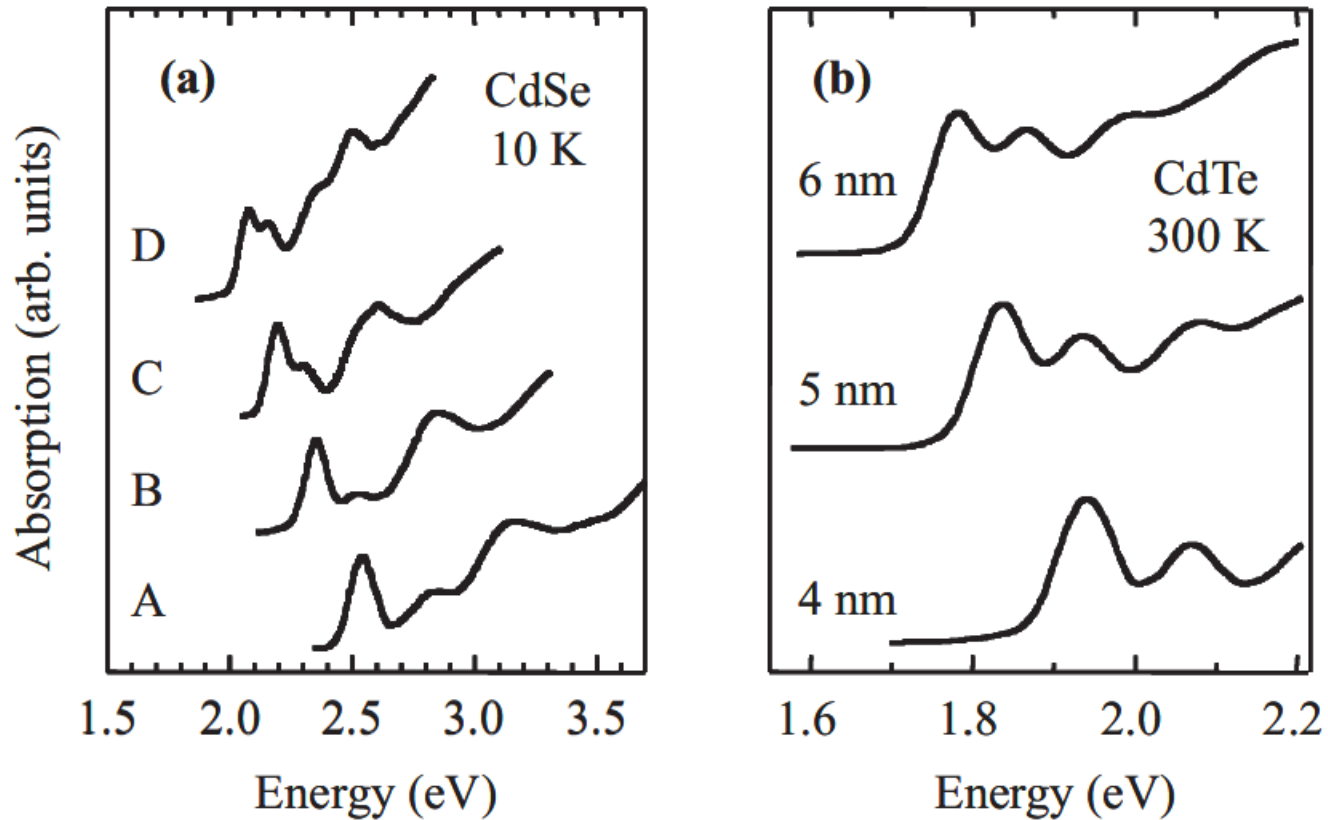
Initial imperfections on the substrate
Shadow effect leads to tilted column growth.

Colloidal quantum dots

Color Control by QD Particle Size



Colloidal quantum dots

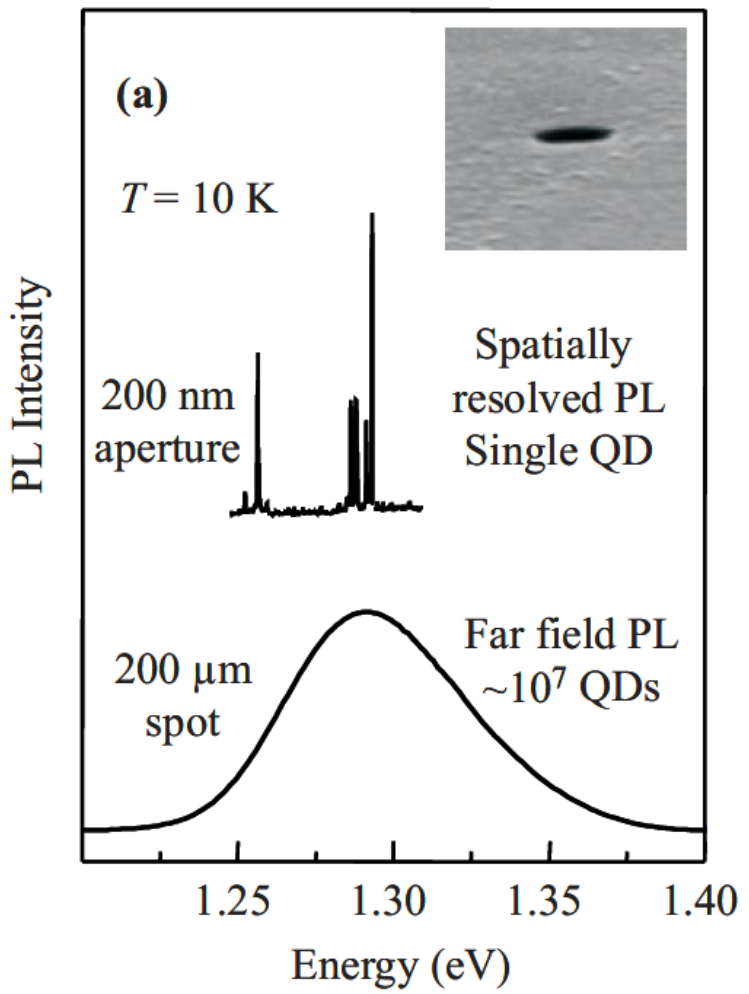
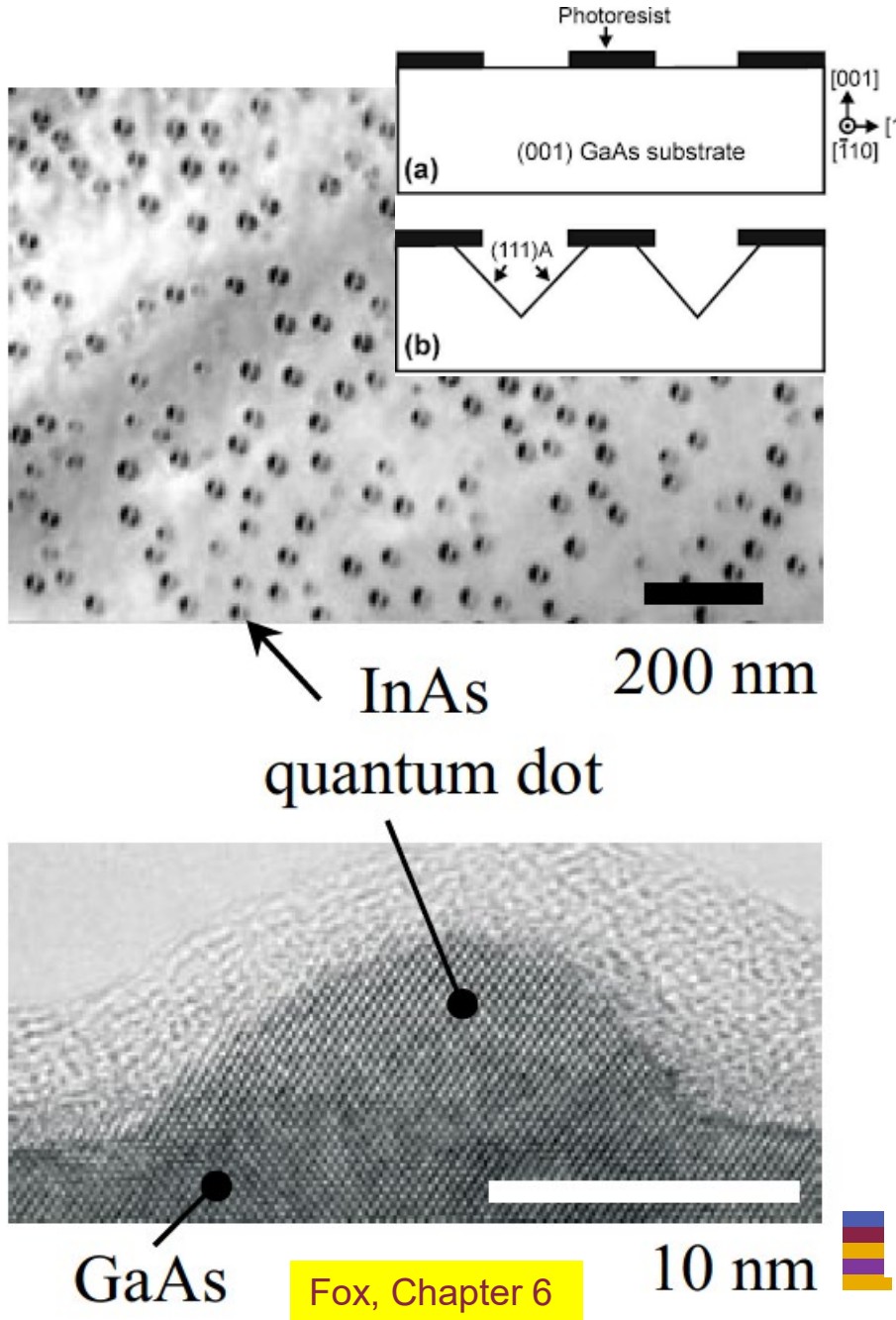


Fox, Chapter 6

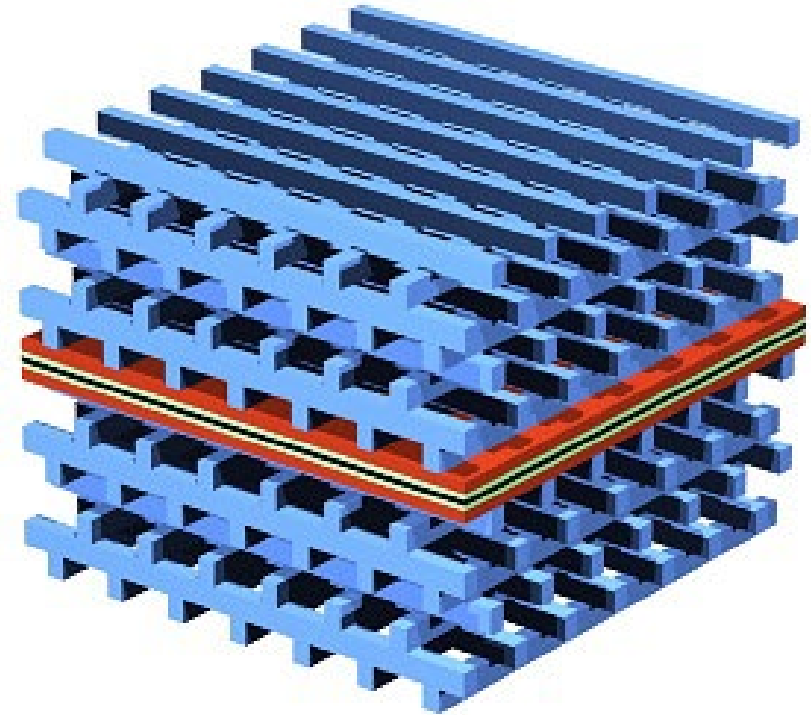
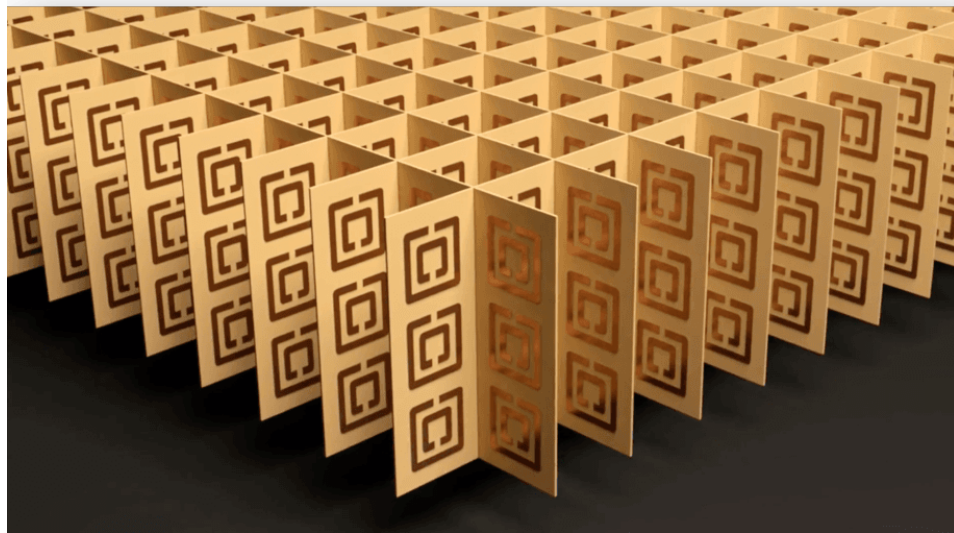
C.R. Kagan, C.B. Murray, M.G. Bawendi, PRB **54**, 8633 (1996)

Epitaxial quantum dots

Stranski-Krastanov growth of semiconductor islands by MBE



Metamaterials



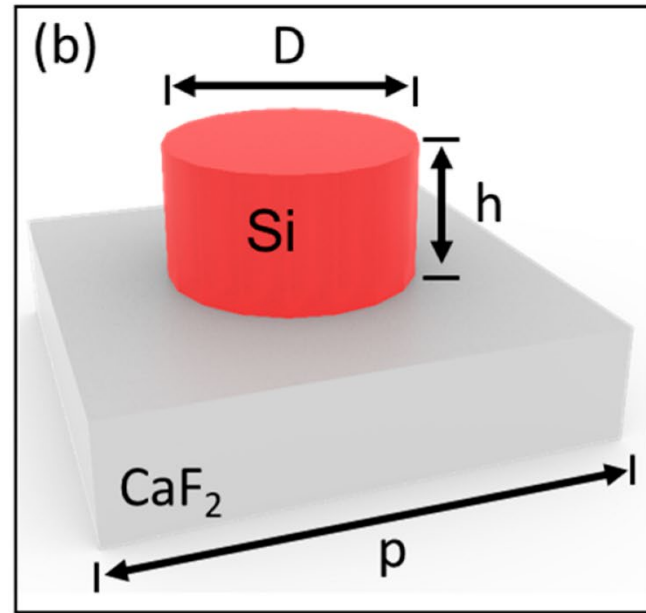
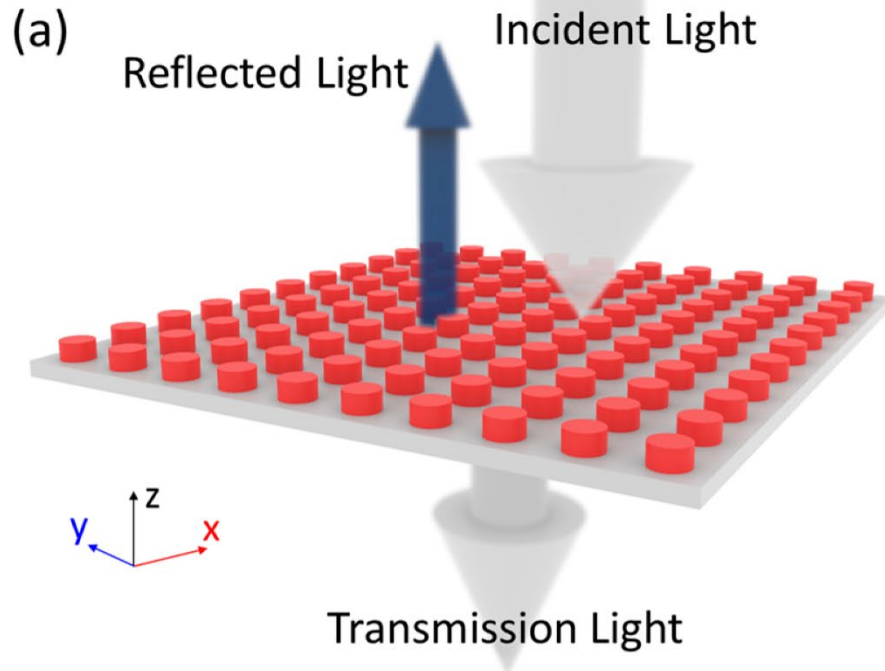
Faraday: $\vec{k} \times \vec{E}_0 = \omega \mu \mu_0 \vec{H}_0$

Metamaterials are artificially structured three-dimensional materials with feature size less than a wavelength. They have unusual properties:
Photonic band gaps, left-handed materials, negative index.

Sajeev John
Shelby, APL 78, 489 (2001)



Metasurfaces



Metasurfaces can be designed to have reflection, absorption, and emission properties not possible in homogeneous materials. Especially useful for sensors and antennas.

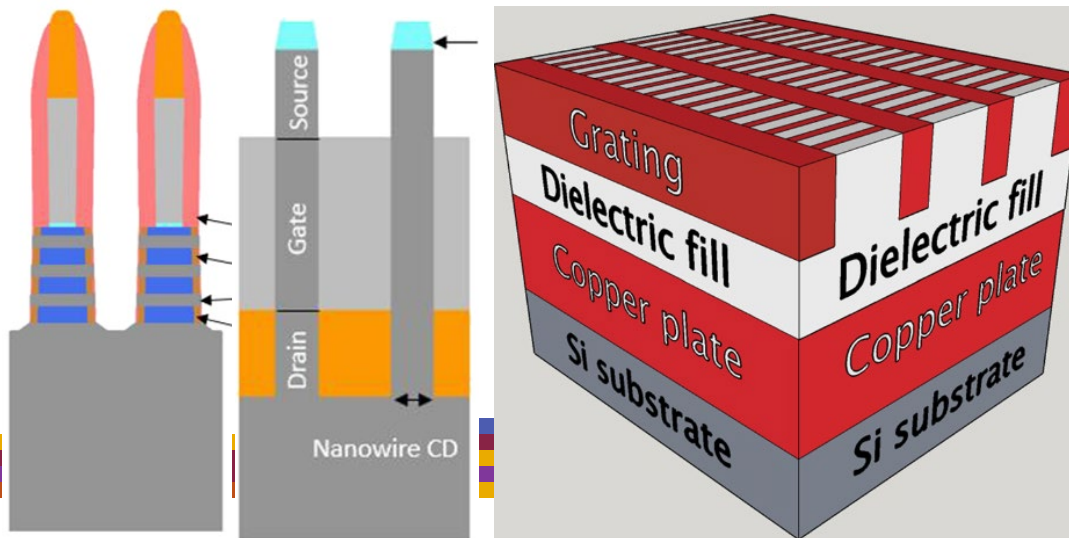
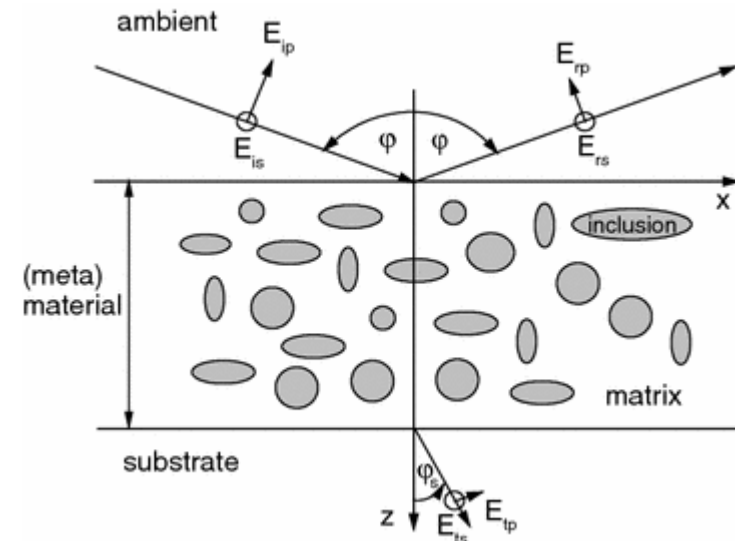
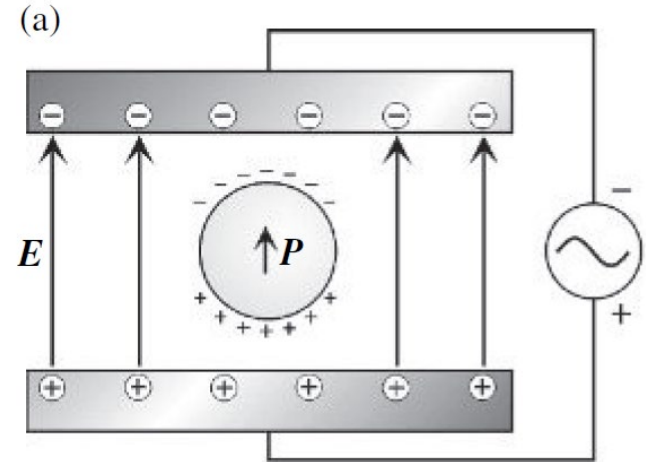
Modeling Metamaterials and Metasurfaces

Bruggeman effective medium approximation (BEMA)

$$\sum_{i=1}^n f_i \frac{\epsilon_i - \epsilon}{\epsilon_i + 2\epsilon} = 0$$

Rigorous coupled wave approximation:

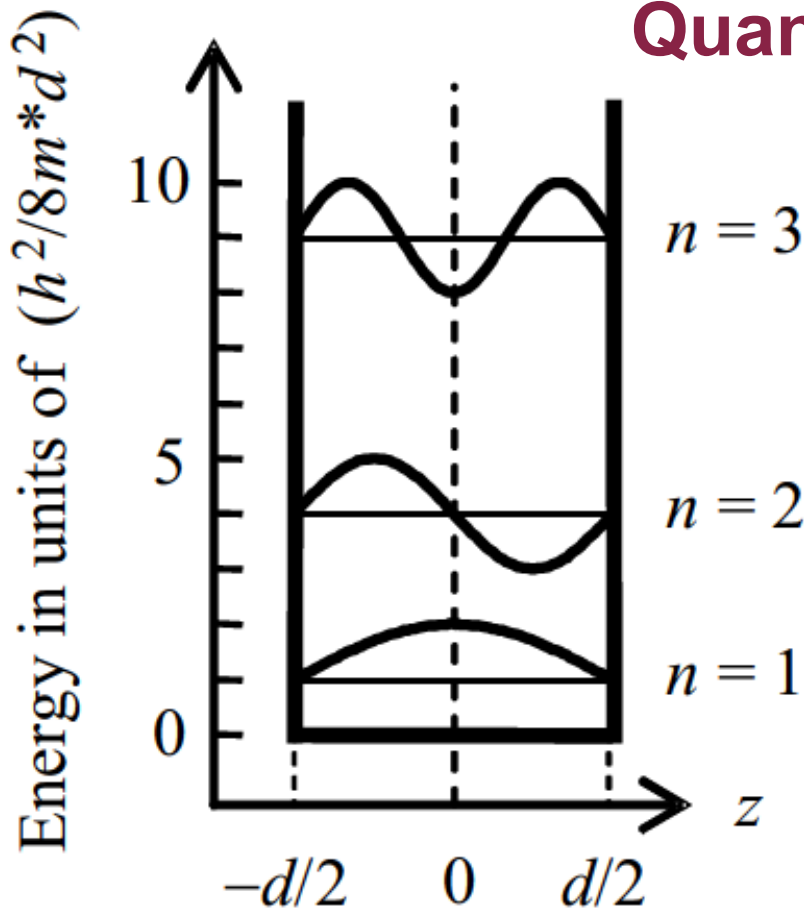
- Fourier method for periodic structures
- Numerical solution of Maxwell's equations



Schmidt, JAP **114083519** (2013)
Diebold, APL Mater. **6**, 058201 (2018)



Quantum well with infinite barriers



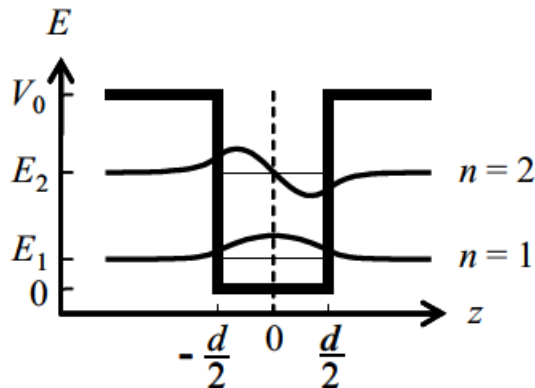
$$\Psi(x, y, z) = \psi(x, y)\varphi(z)$$

$$E(\vec{k}, n) = \frac{\hbar^2 k^2}{2m} + \frac{\hbar^2}{2m} \left(\frac{n\pi}{d}\right)^2$$

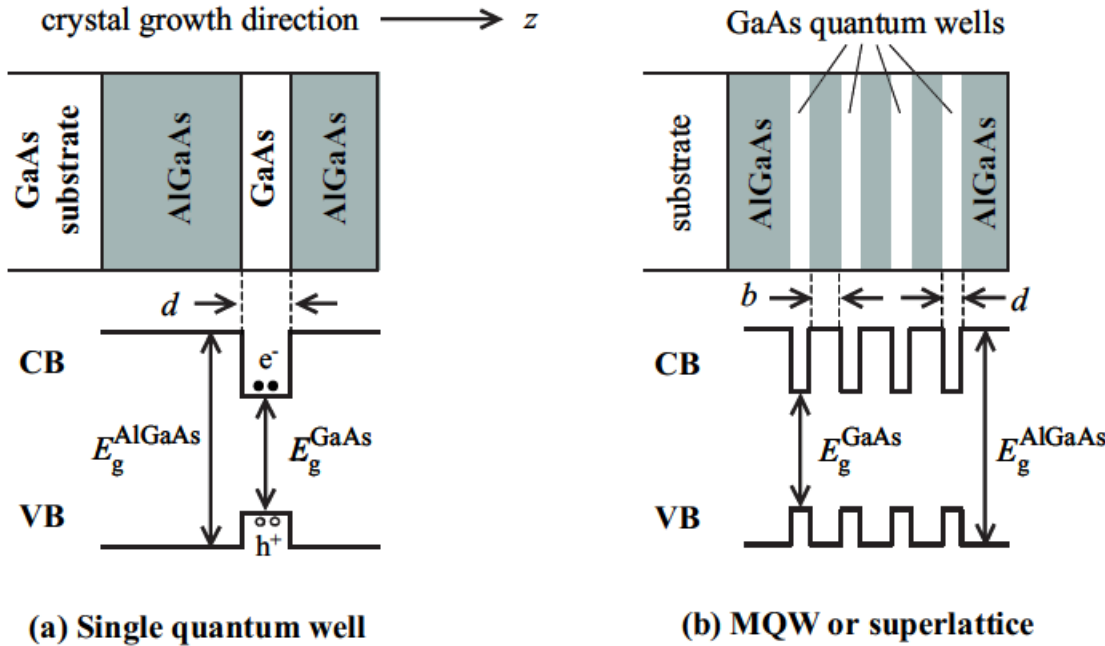
Use carrier effective mass.
(Electrons more affected than holes).

Finite barrier:

Confinement energies are lower.
Wave function leaks into barrier.
Finite number of bound states.
Solve numerically.



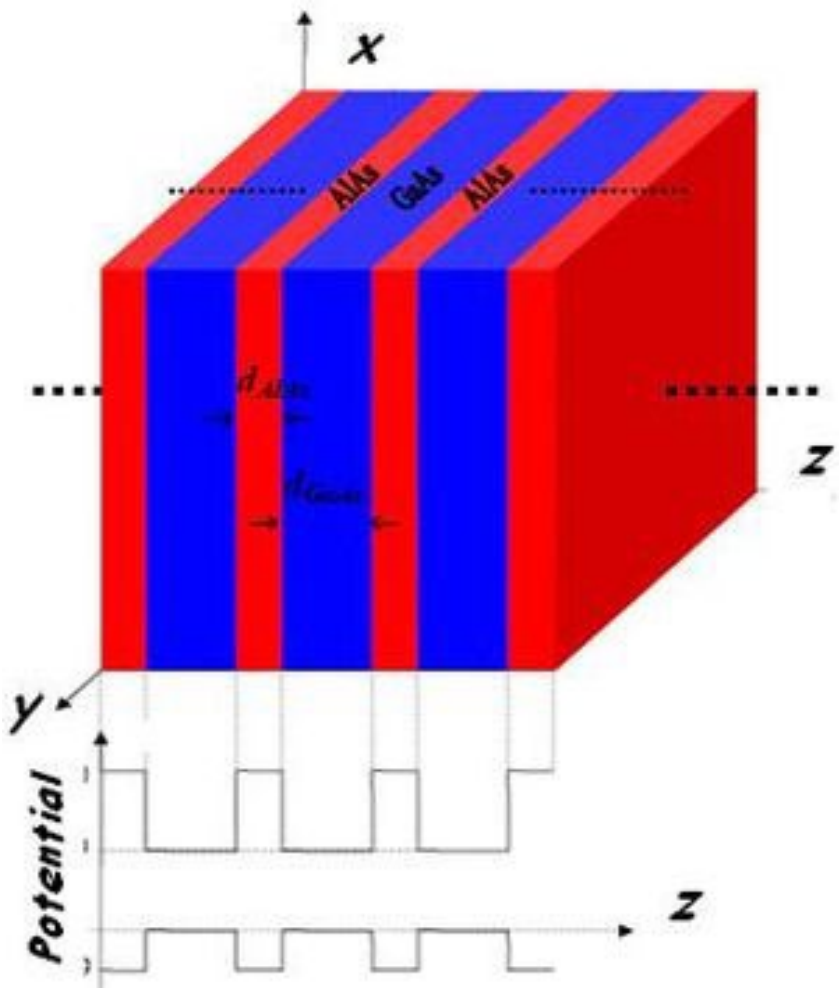
Quantum wells and superlattices (or MQW)



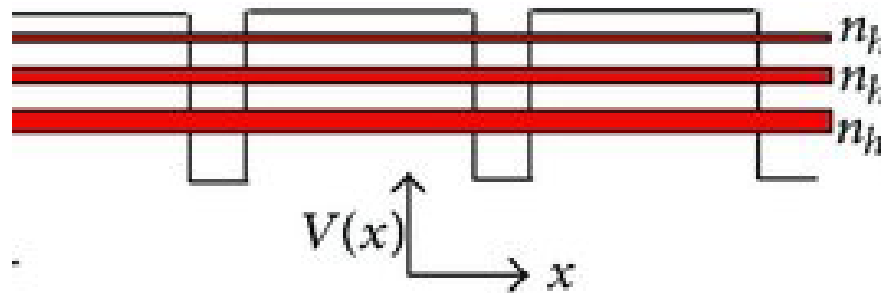
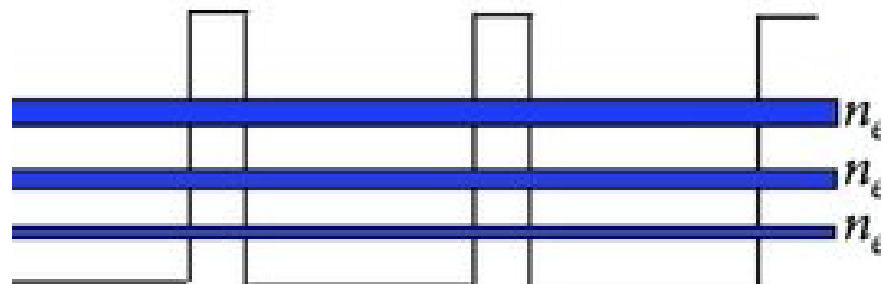
Quantum well:
Superlattice:

wave function entirely contained in one well.
wave function leaks through the barrier into the next well (barrier thin or low).

Superlattice minibands



CB

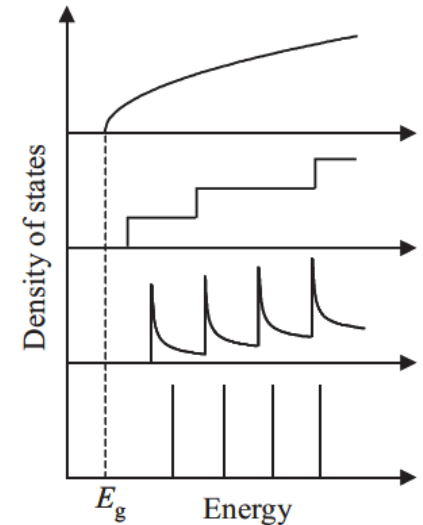
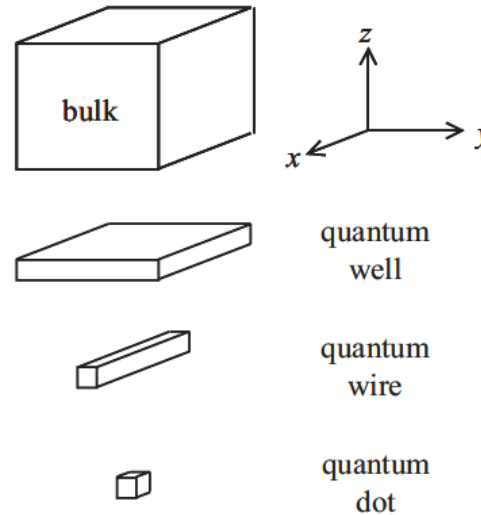
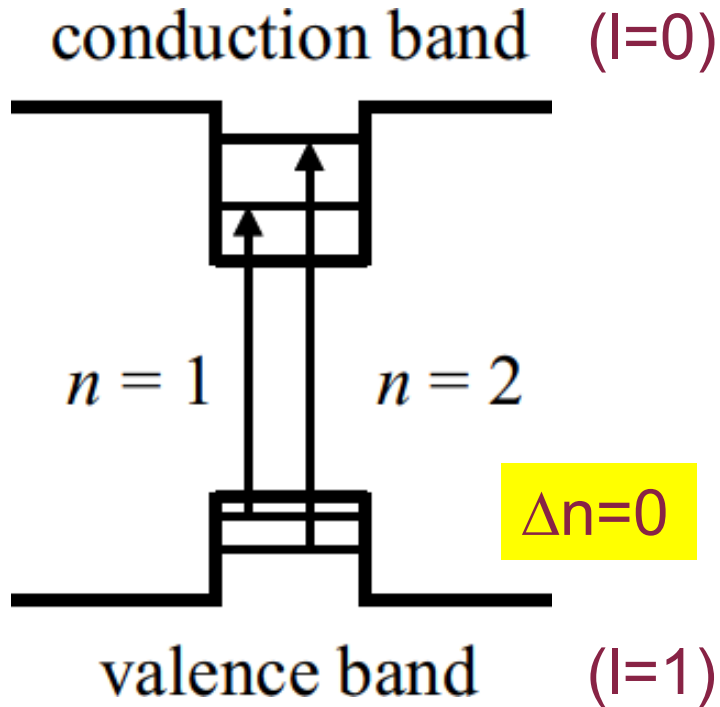


VB

Enhanced absorption/emission in quantum structures

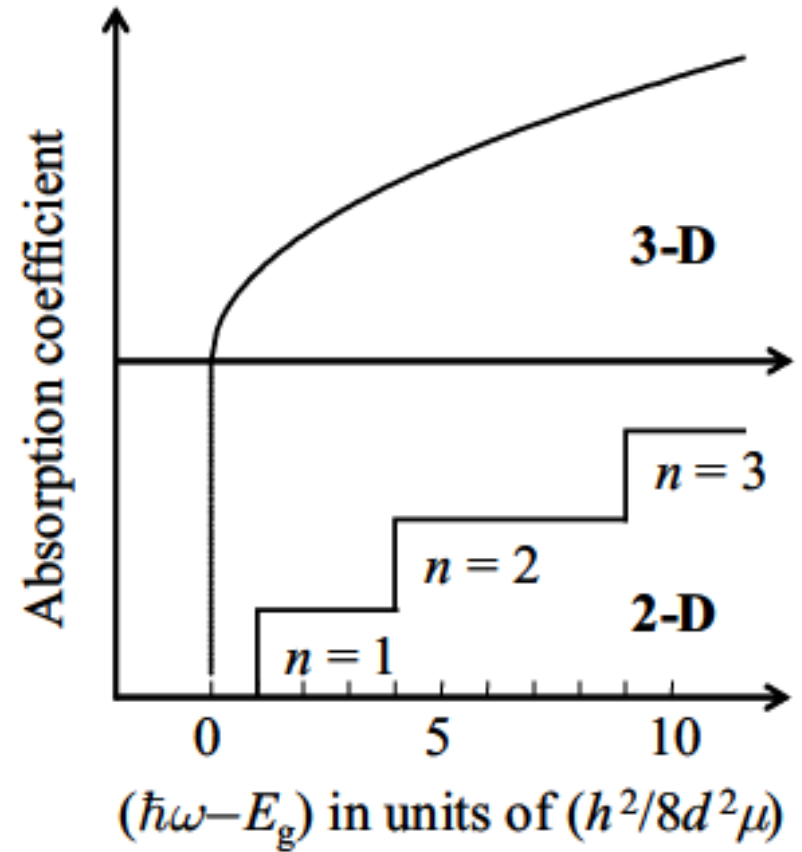
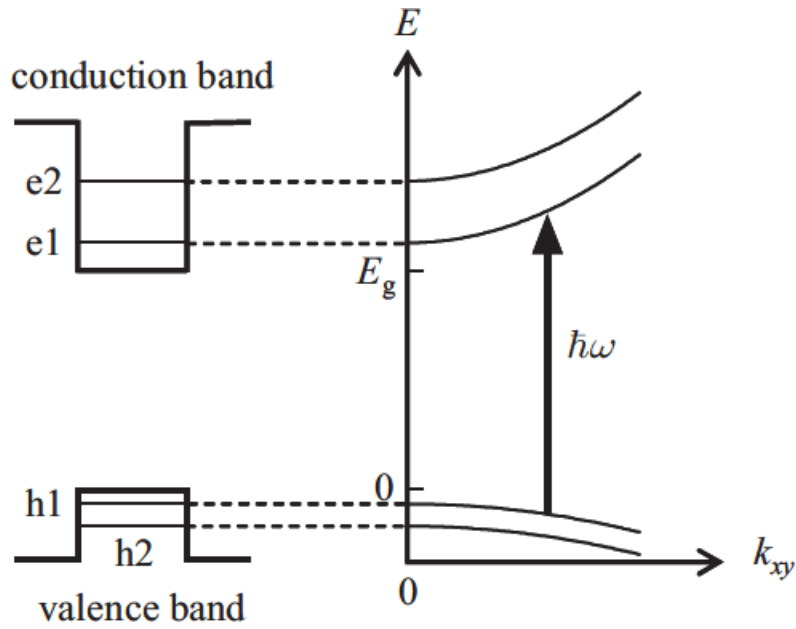
$$\frac{1}{\tau} = \frac{2\pi}{\hbar} |\langle f | H_{eR} | i \rangle|^2 g_{fi}(\hbar\omega)$$

Confinement enhances overlap matrix element



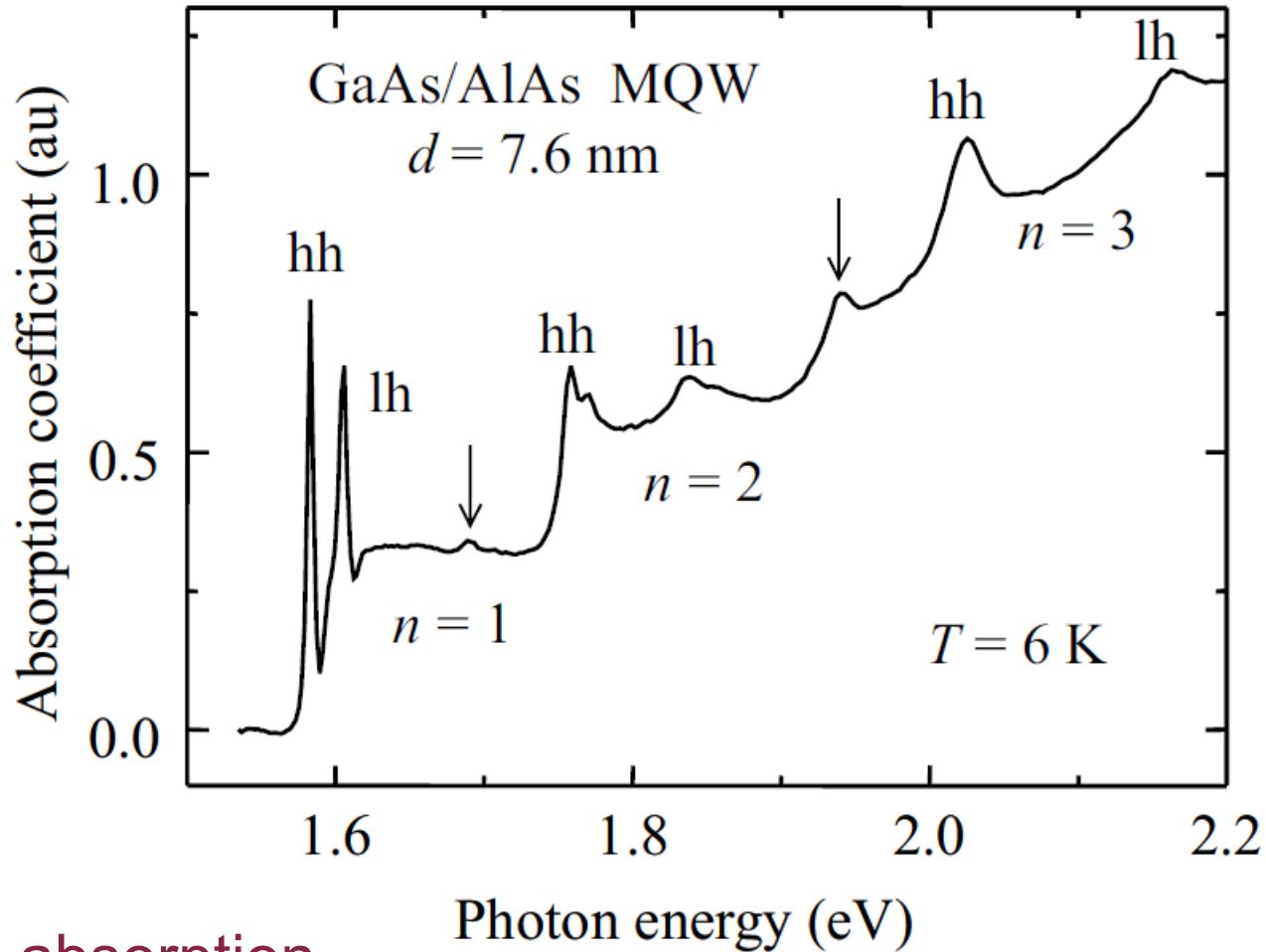
Modified density of states

Quantum well absorption



Electron and hole subbands

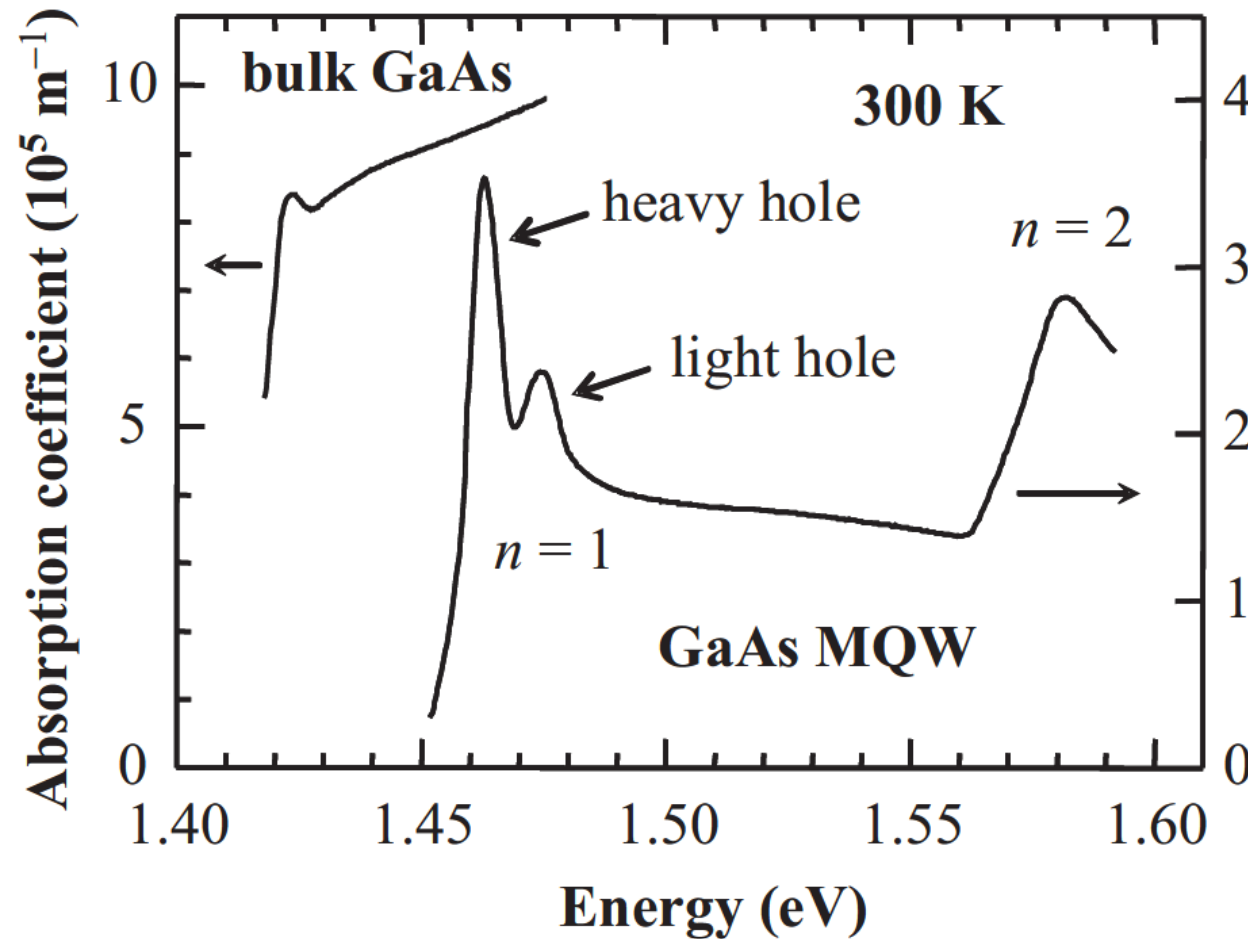
Quantum well absorption in GaAs/AlAs MQW



Steps: 2-D absorption

Peaks: Discrete excitons (arrows show forbidden transitions)

Excitonic effects are enhanced in quantum wells



Bulk GaAs:
 $R_x = 4.2 \text{ meV}$

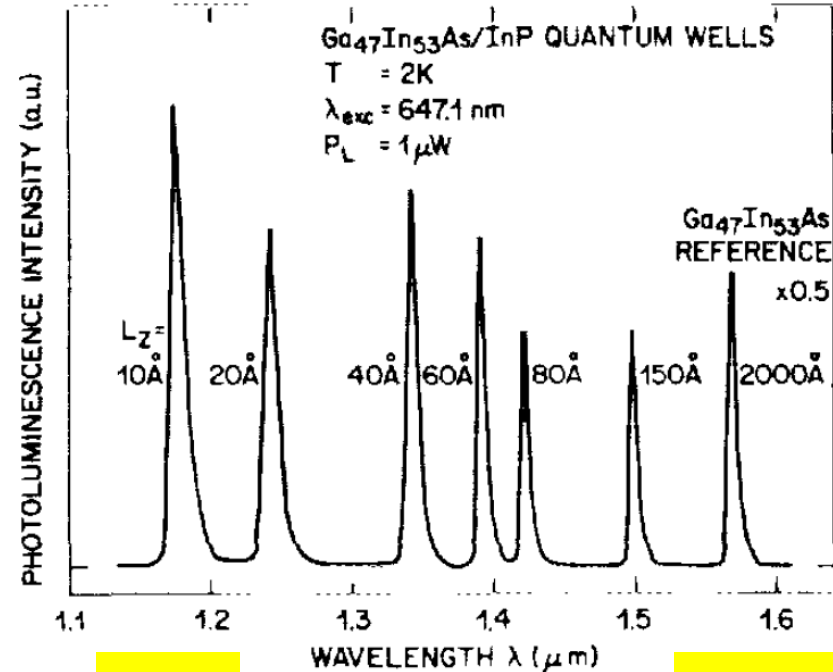
Here ($d=10 \text{ nm}$):
 $R_x = 10 \text{ meV}$
 Increased overlap
 integral

hh/lh splitting:
 Confinement or
 strain ???

$$E(\vec{k}, n) = \frac{\hbar^2 k^2}{2m} + \frac{\hbar^2}{2m} \left(\frac{n\pi}{d} \right)^2$$

Confinement shifts in quantum wells

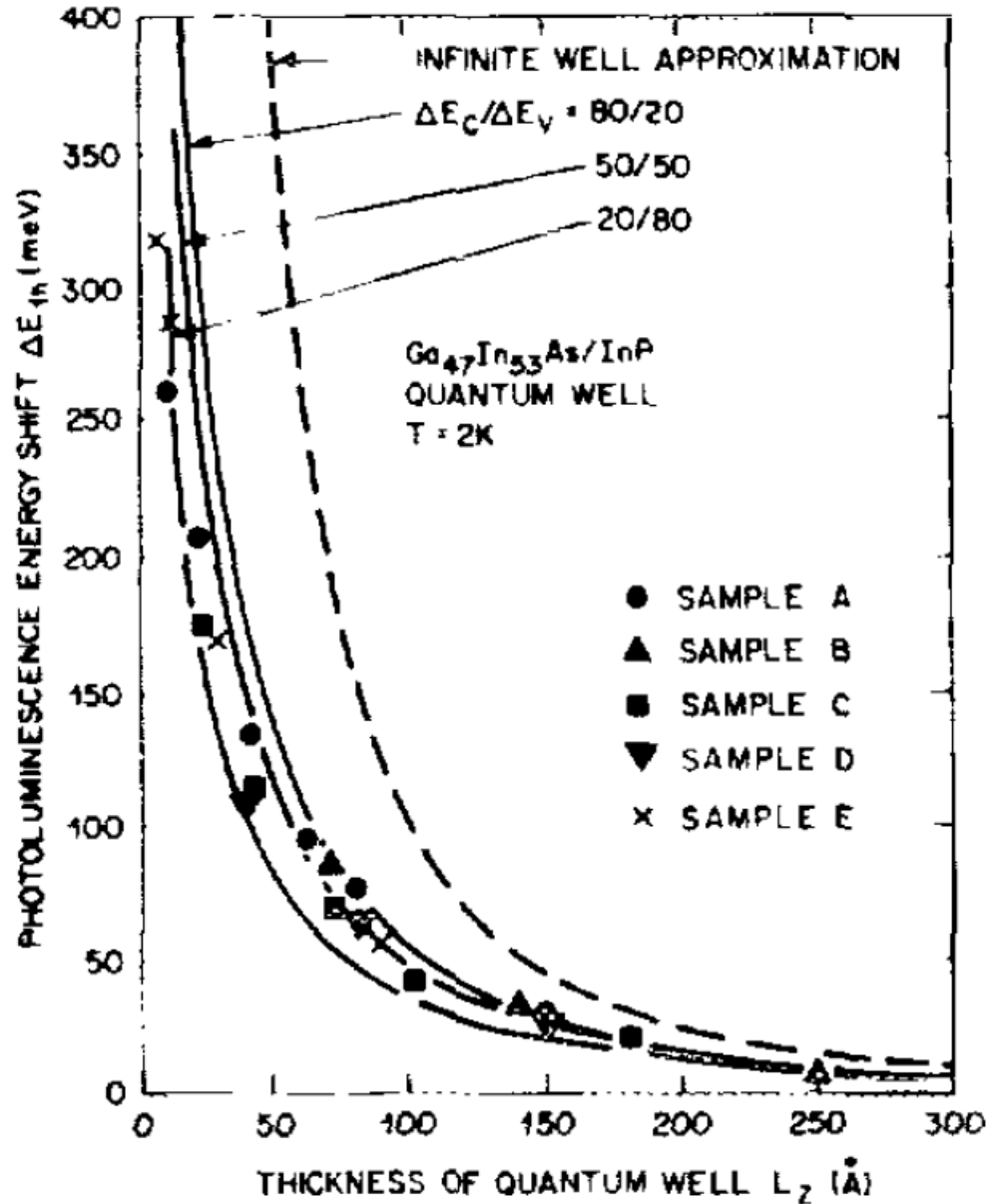
$$\hbar\omega = E_g + E_{hh1} + E_{e1}$$



thin

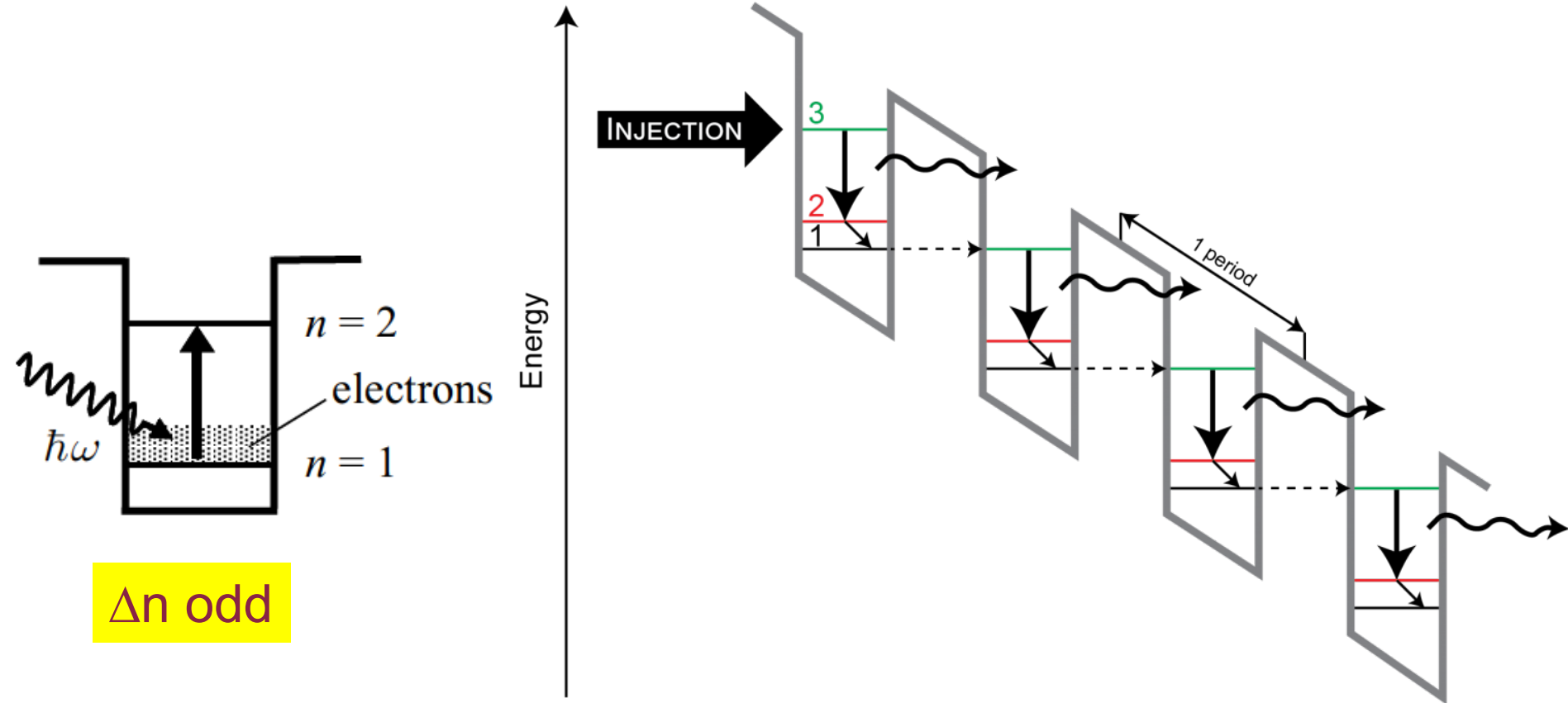
thick

$$\Delta E = \frac{\hbar^2}{2m} \left(\frac{n\pi}{d} \right)^2$$



W.T. Tsang & E.F. Schubert,
Appl. Phys. Lett. **49**, 220 (1986)

Intersubband transitions, quantum cascade lasers



Infrared detectors and lasers

Light polarized along z (beam in quantum well plane)



Defects

Vacancy

V_A

missing atom

Interstitial

I_A

extra atom between lattice sites

Substitutional

C_A

defect atom replaces host atom

Antisite

atoms are switched

Frenkel defect pair

V_A-I_A

atom moved to interstitial site

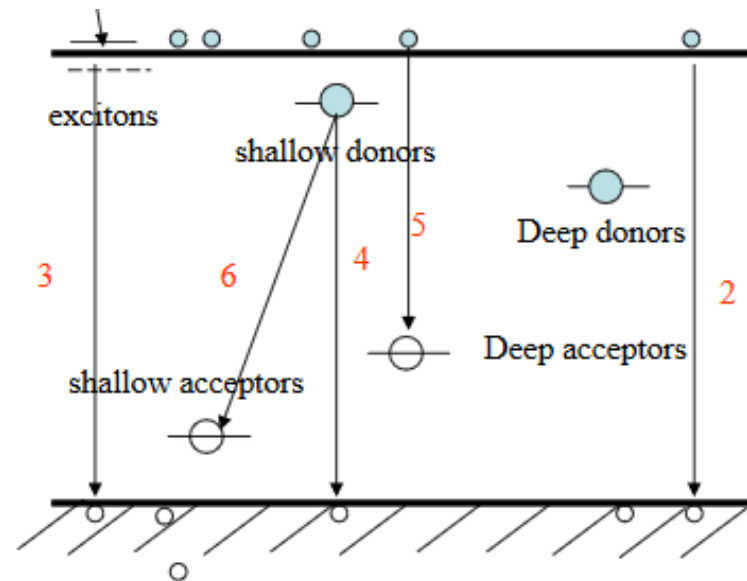
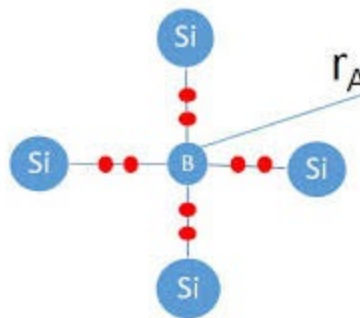
More complex defects (combinations)

Donor: Substitutional defect, adds an electron

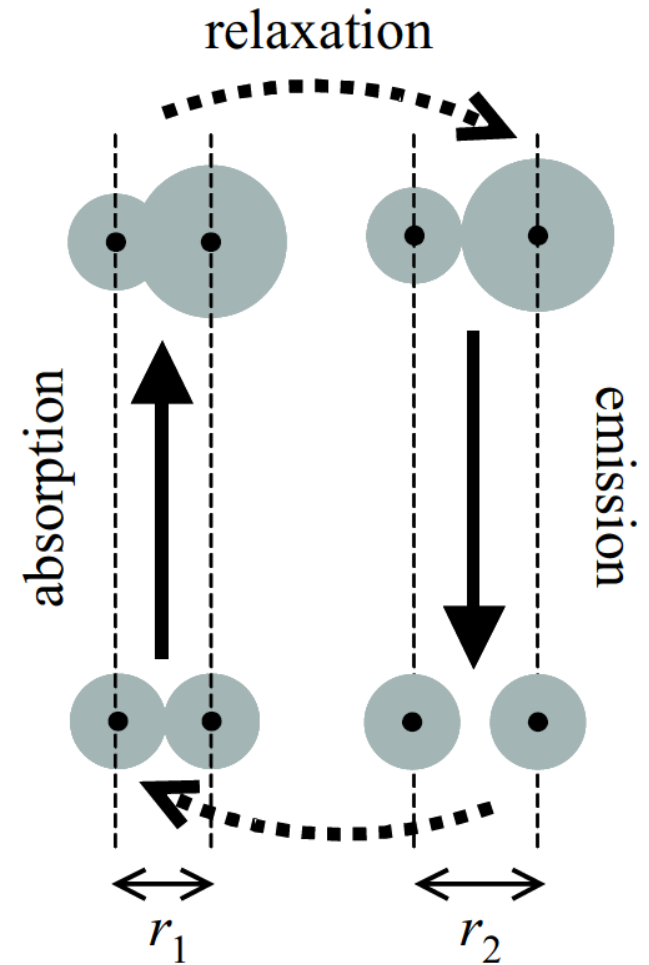
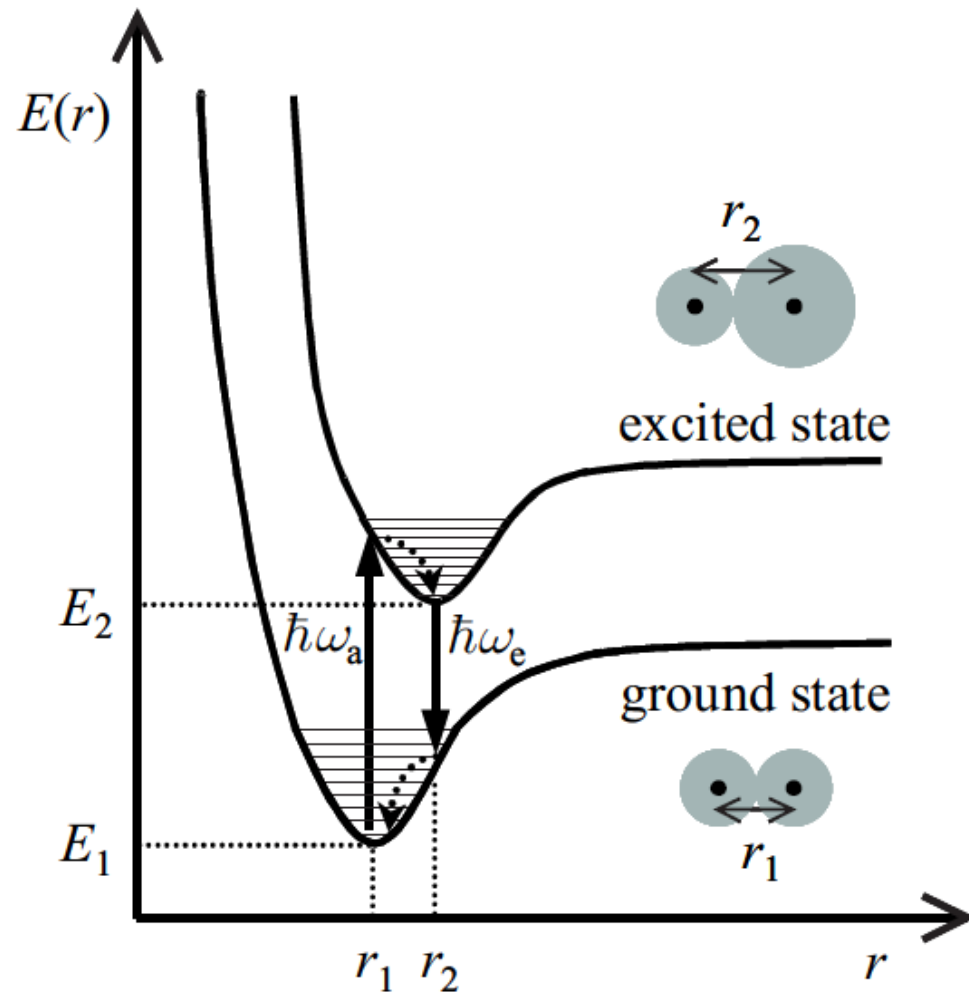
Acceptor: Substitutional defect, adds a hole

Isoelectronic: Impurity has same number of electrons as host

Shallow and **deep** defects.



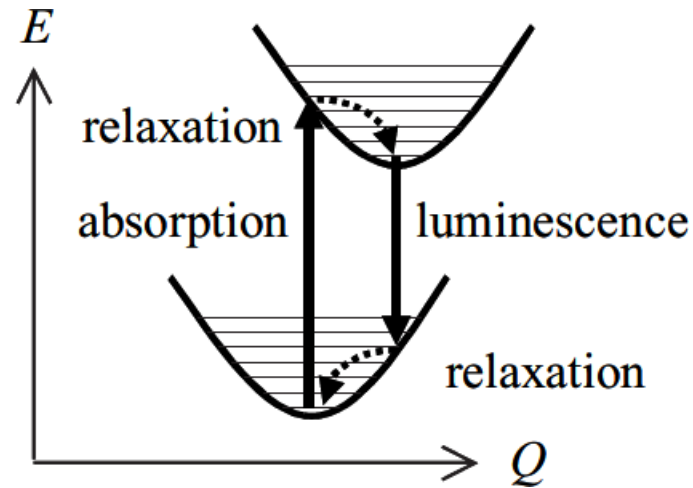
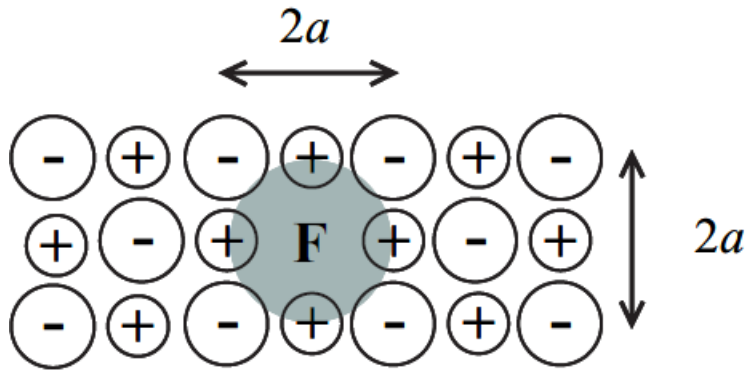
Defects: Frank-Condon principle



Born-Oppenheimer approximation

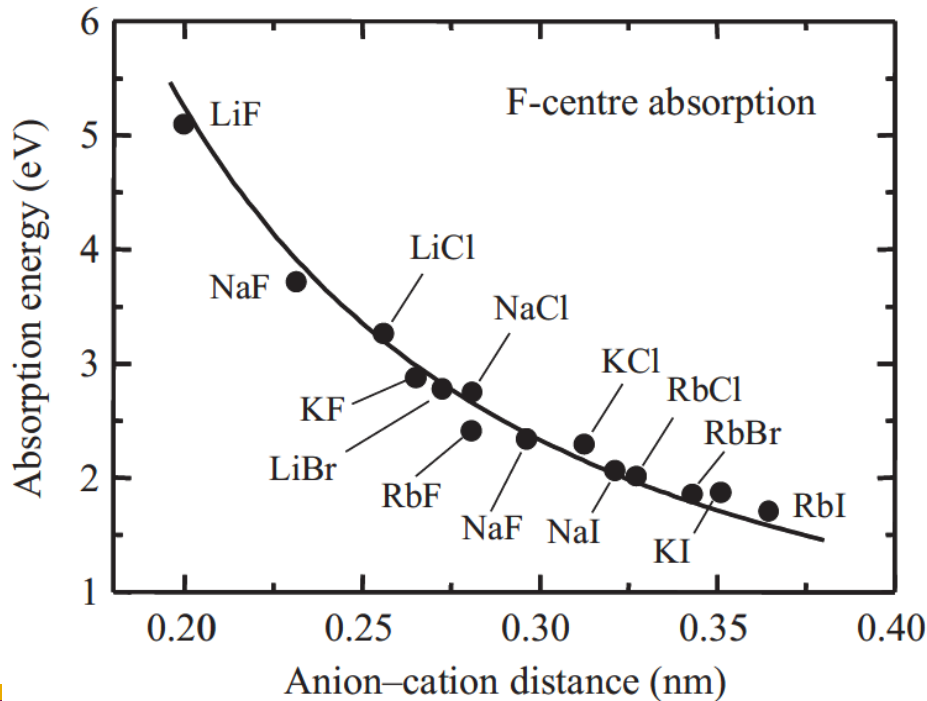


Defects: Color centers (F-centers)

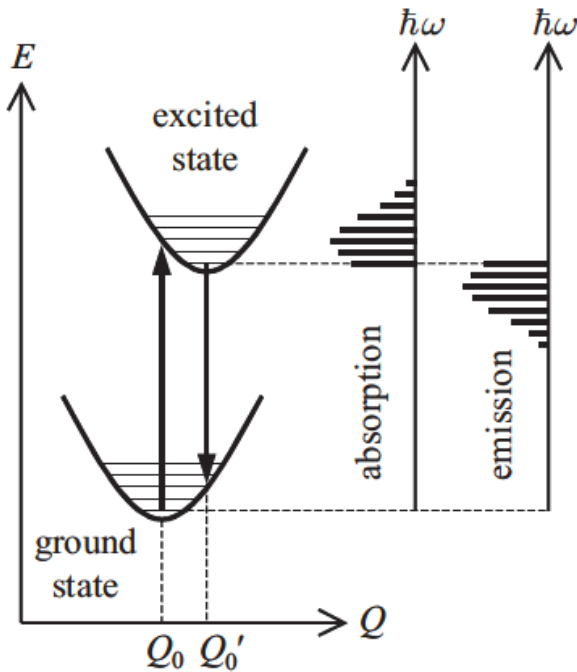
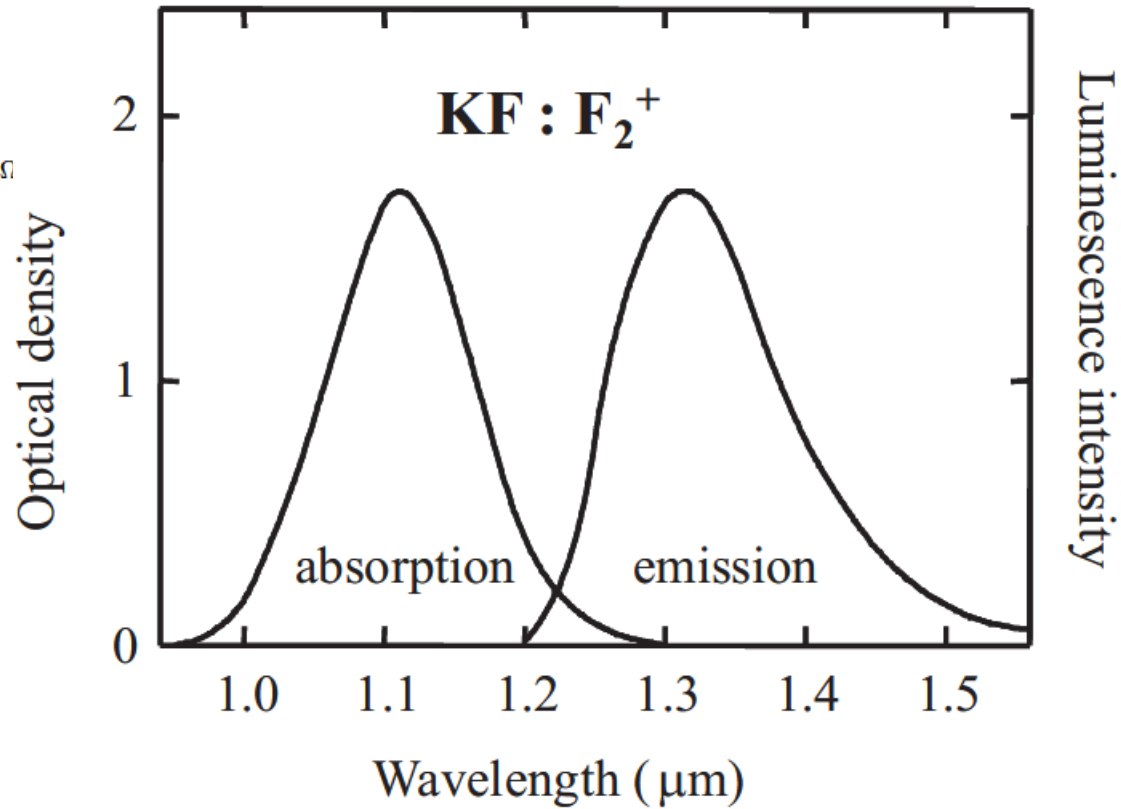
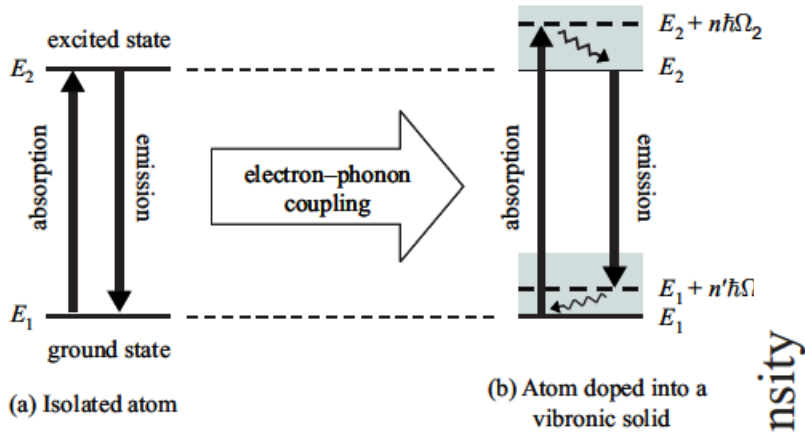


Alkali halide: large-gap insulator
 Vacancy: positively charged hole

F-center: electron trapped at anion vacancy.



Defects: Shift between absorption and emission



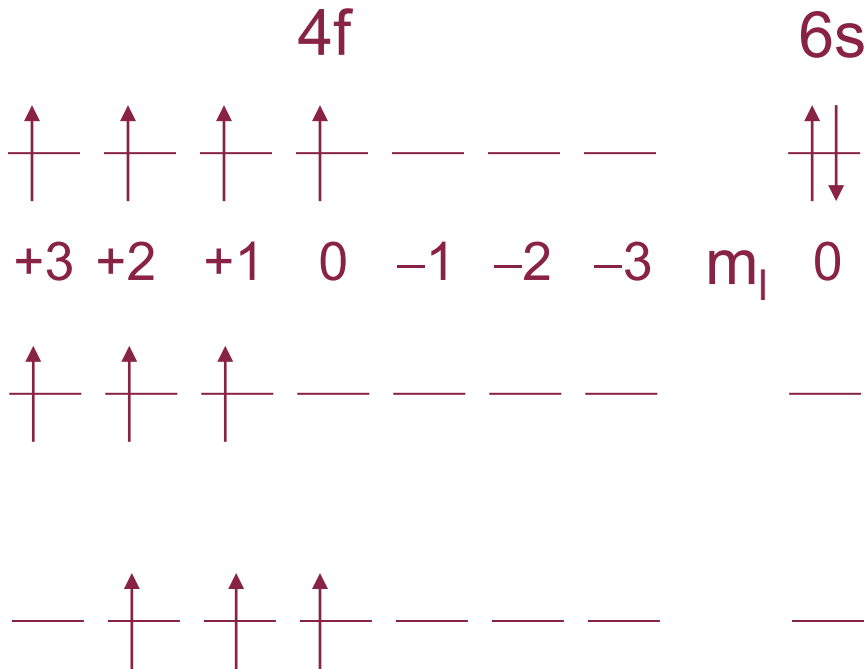
Two vacancies, one electron
Color-center lasers

Hund's Rules (applied to Neodymium, Z=60)

1. Maximize S
2. Maximize L
3. If shell is less than half full, $J=L-S$
If shell is more than half full, $J=L+S$.

$$2S+1L_J$$

Russell-Saunders
(LS) coupling



atom: $5I_4$

Nd: $4f^4 6s^2 S=2, L=6, J=4$

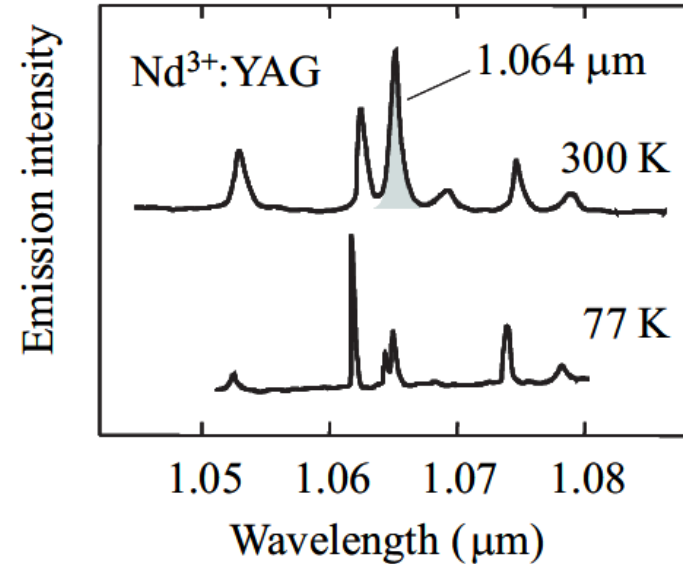
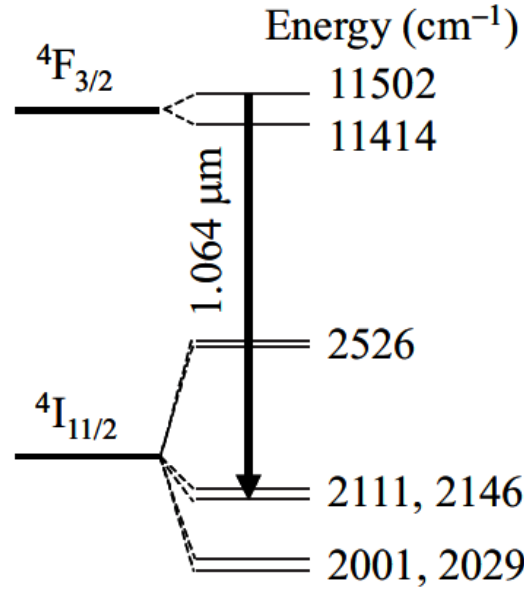
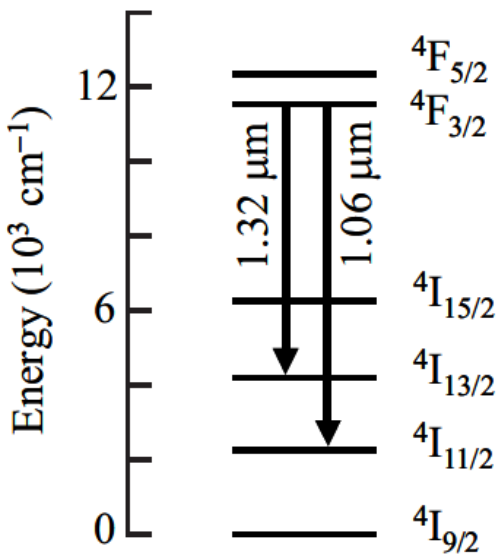
3+ ion ground state: $4I_{9/2}$

Nd³⁺: $4f^3 6s^0 S=3/2, L=6, J=9/2$

3+ ion excited state: $4F_{3/2}$

Nd³⁺: $4f^3 6s^0 S=3/2, L=3, J=3/2$

Defects: Rare earth metal ions in insulator (Nd:YAG)



Atomic configuration

Weak crystal-field splitting

$\Delta J=4$ or 5 forbidden

Transition weakly allowed with crystal-field splitting

YAG: Y₃Al₅O₁₂

Nd:YAG lasers

YAG: 1064 nm

YLiF₄: 1053 nm

Broad (120 GHz)



Defects: Transition metal ions (Ti:sapphire)

3d

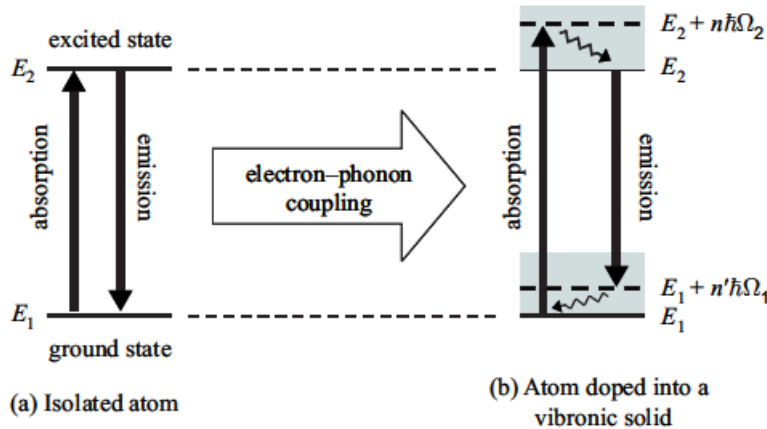
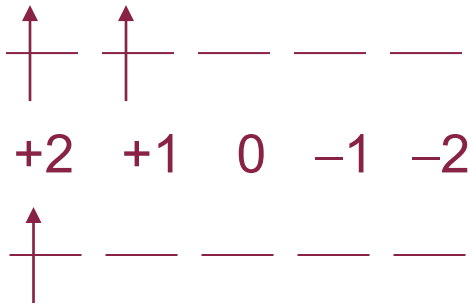
4s

atom: 3F_2

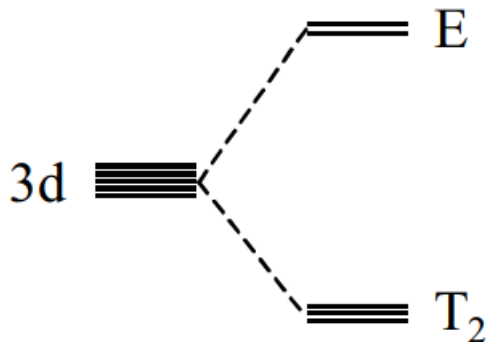
Ti: $3d^2 4s^2$ $S=1, L=3, J=2$

ion: ${}^2D_{3/2}$

Ti³⁺: $3d^1 4s^0$ $S=1/2, L=2, J=3/2$



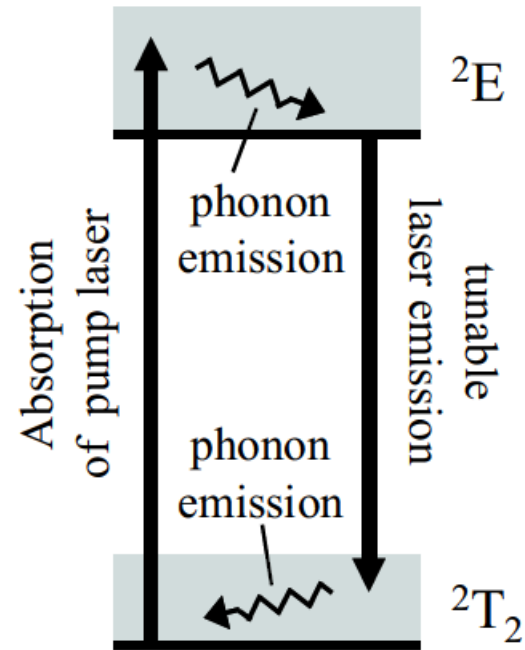
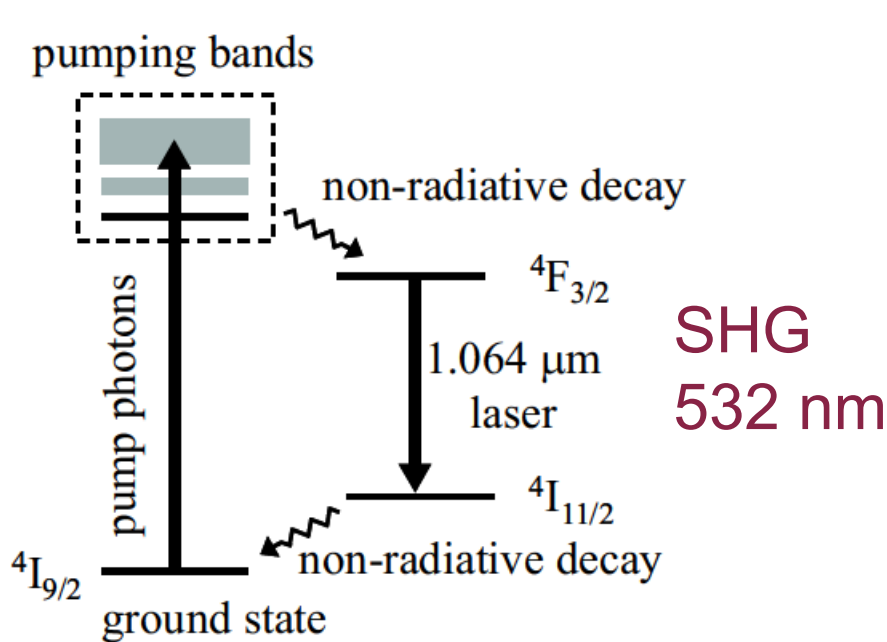
Strong vibronic coupling
Broad lines (gain spectrum)
Good for ultrafast lasers



Strong crystal-field splitting
Depends on crystal environment

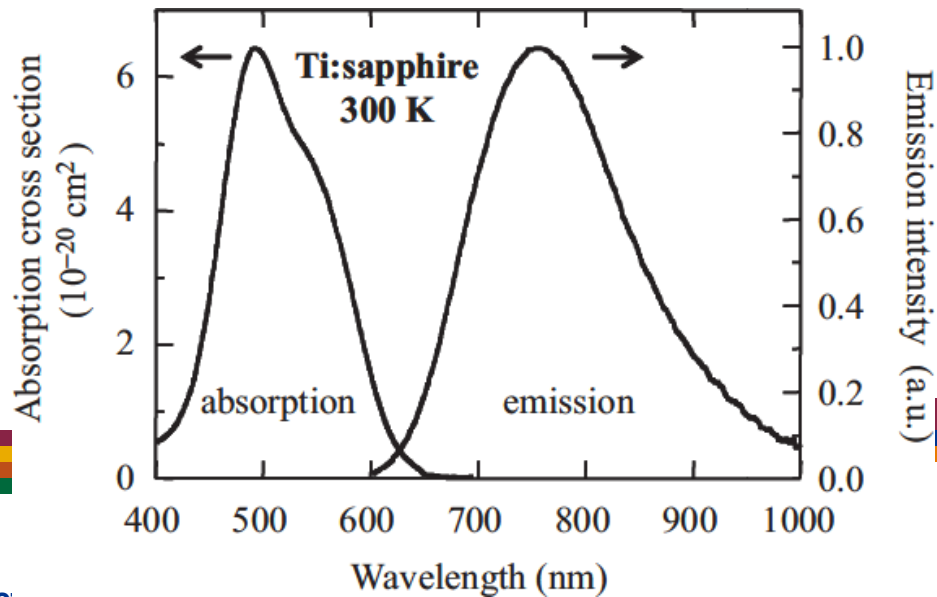


Solid-State Lasers: Ti:sapphire and Nd:YAG

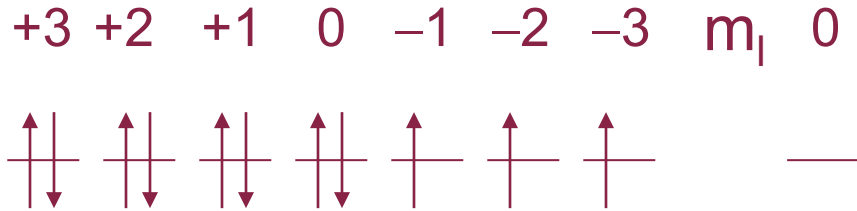


Pumping with

- Flash lamp (pulsed)
- Xe lamp (CW)
- GaAs diode (efficient)

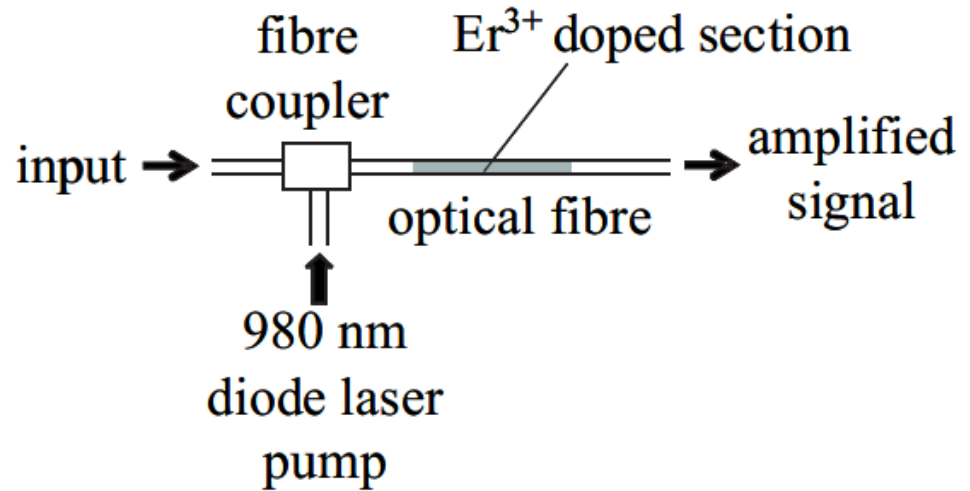
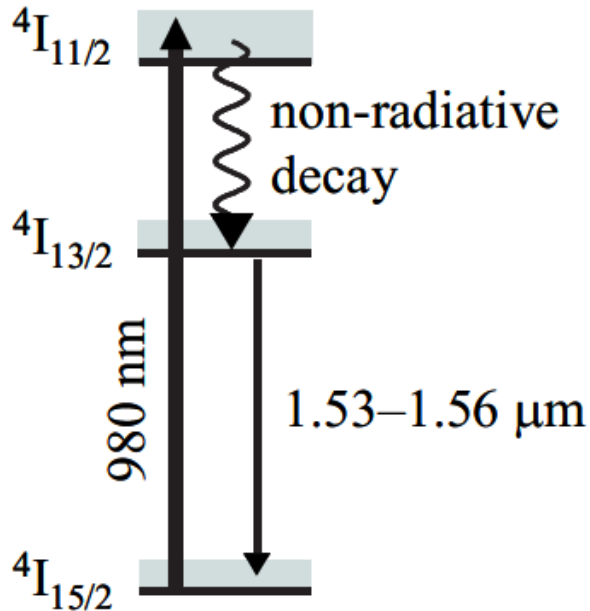


Erbium fiber laser



3+ ion ground state: $4I_{15/2}$

Er³⁺: $4f^{11} 6s^0$ S=3/2, L=6, J=15/2



Pumped with 980 nm diode laser.
Important for telecommunications.

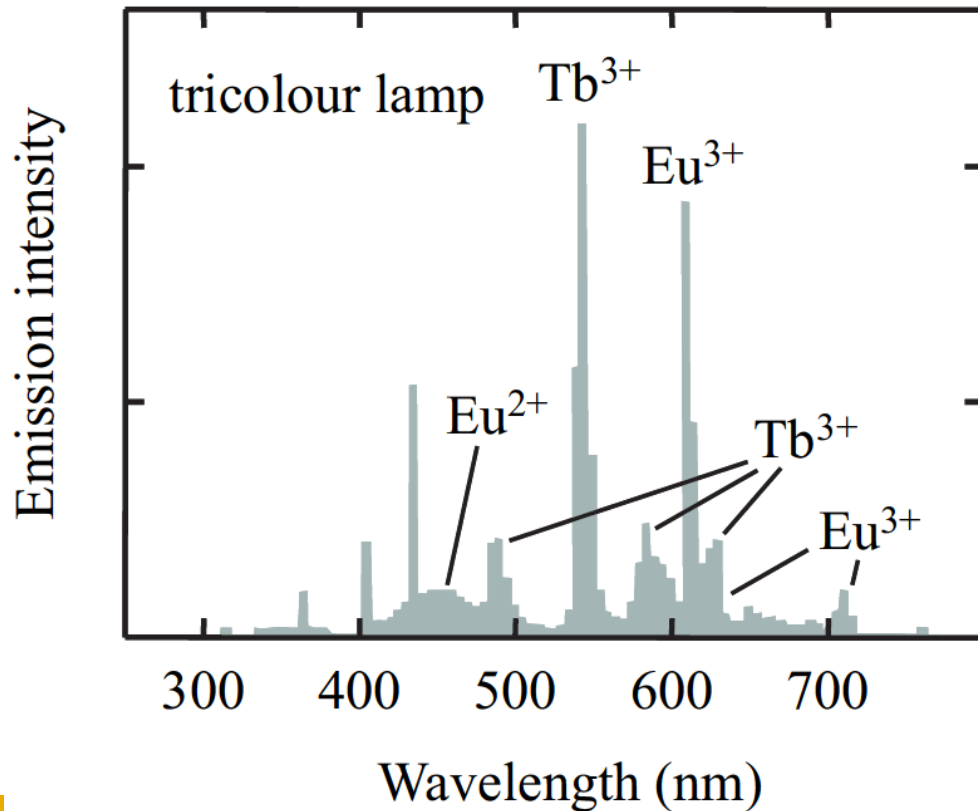


Phosphors

Obsolete: Cathode ray tubes (TV, oscilloscope)

Fluorescent tubes.

Convert discrete LED (or Hg) emission into white light.

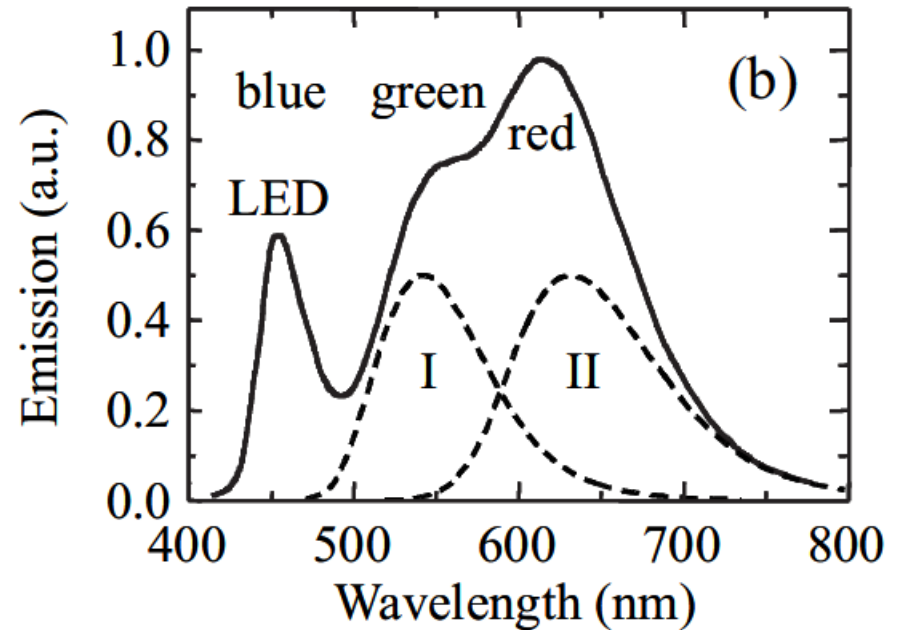
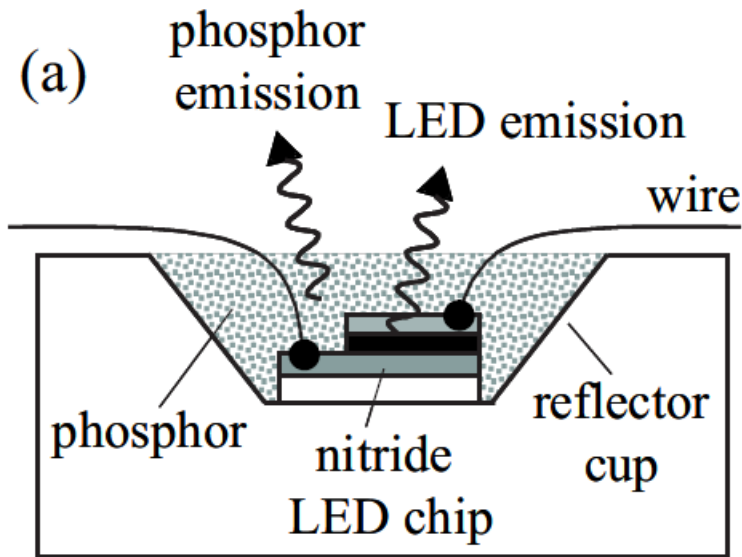


$\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu}^{2+}$ ($4f^7$)
450 nm (blue)

$\text{CeMgAl}_{11}\text{O}_{19}:\text{Tb}^{3+}$ ($4f^8$)
550 nm (green)

$\text{Y}_2\text{O}_3:\text{Eu}^{3+}$ ($4f^6$)
610 nm (red)

White light emitting diodes (white LEDs)



Blue InGaN LED with green and red phosphors

I: $\text{SrSi}_2\text{O}_2\text{N}_2:\text{Eu}^{2+}$

II: $\text{Sr}_2\text{Si}_5\text{N}_8:\text{Eu}^{2+}$

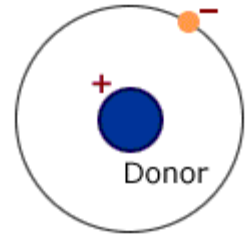
Color temperature: 3200 K

Shallow (hydrogenic) defect: Si donor in GaAs

Extra valence electron

P nucleus positively charged (but screened in crystal)

Hydrogen-like energy spectrum (screened, heavy)



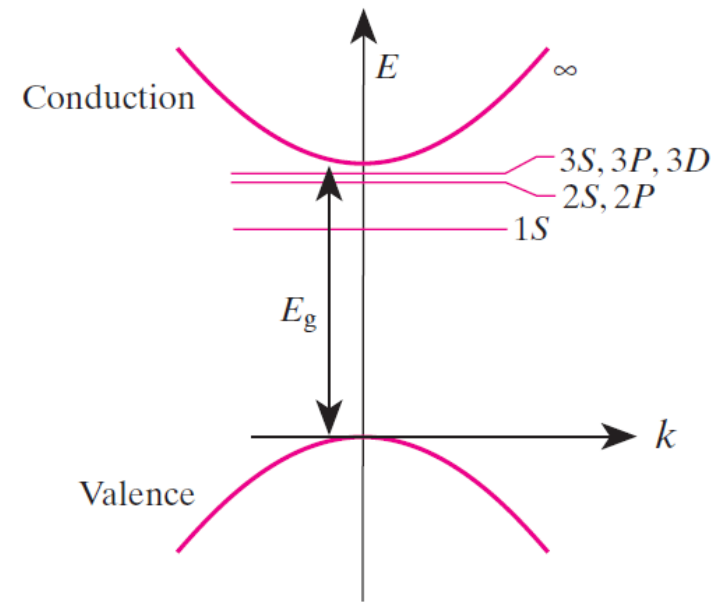
Extra potential: $V_S = + \frac{|e|}{4\pi\epsilon_0 r}$

Rydberg series: $E = E_{CBM} - \frac{R}{n^2}$

Binding energy: $R = \frac{m^*}{m_0} \frac{1}{\epsilon_S^2} \frac{e^4 m_0}{2\hbar^2 (4\pi\epsilon_0)^2}$

Bohr radius.

Similar to exciton problem.



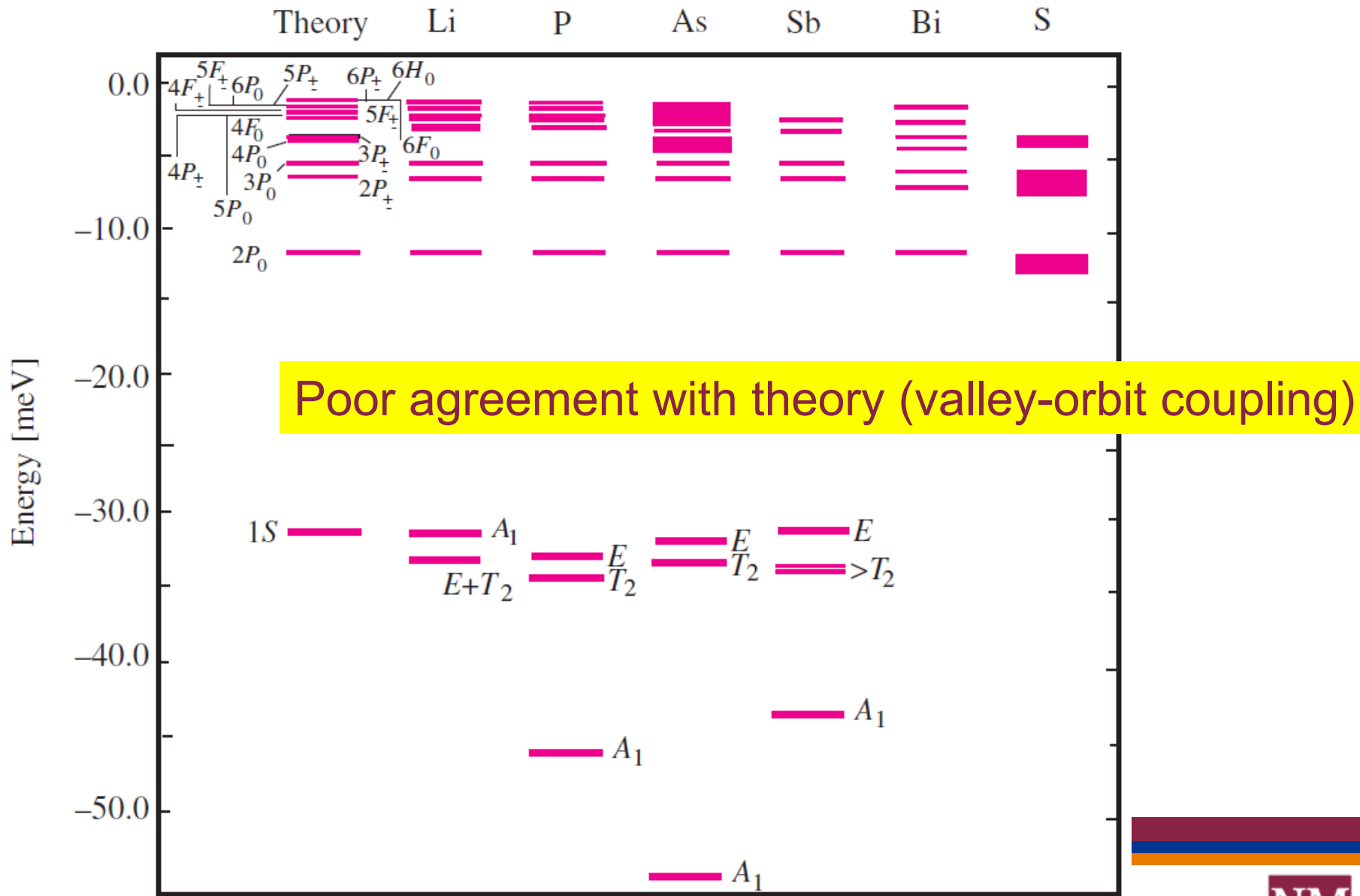
Shallow (hydrogenic) defect: Si donor in GaAs

| Semiconductor | Binding energy from (4.24) [meV] | Experimental binding energy of common donors [meV] |
|---------------|----------------------------------|---|
| GaAs | 5.72 | Si _{Ga} (5.84); Ge _{Ga} (5.88) S _{As} (5.87); Se _{As} (5.79) |
| InP | 7.14 | 7.14 |
| InSb | 0.6 | Te _{Sb} (0.6) |
| CdTe | 11.6 | In _{Cd} (14); Al _{Cd} (14) |
| ZnSe | 25.7 | Al _{Zn} (26.3); Ga _{Zn} (27.9) F _{Se} (29.3); Cl _{Se} (26.9) |

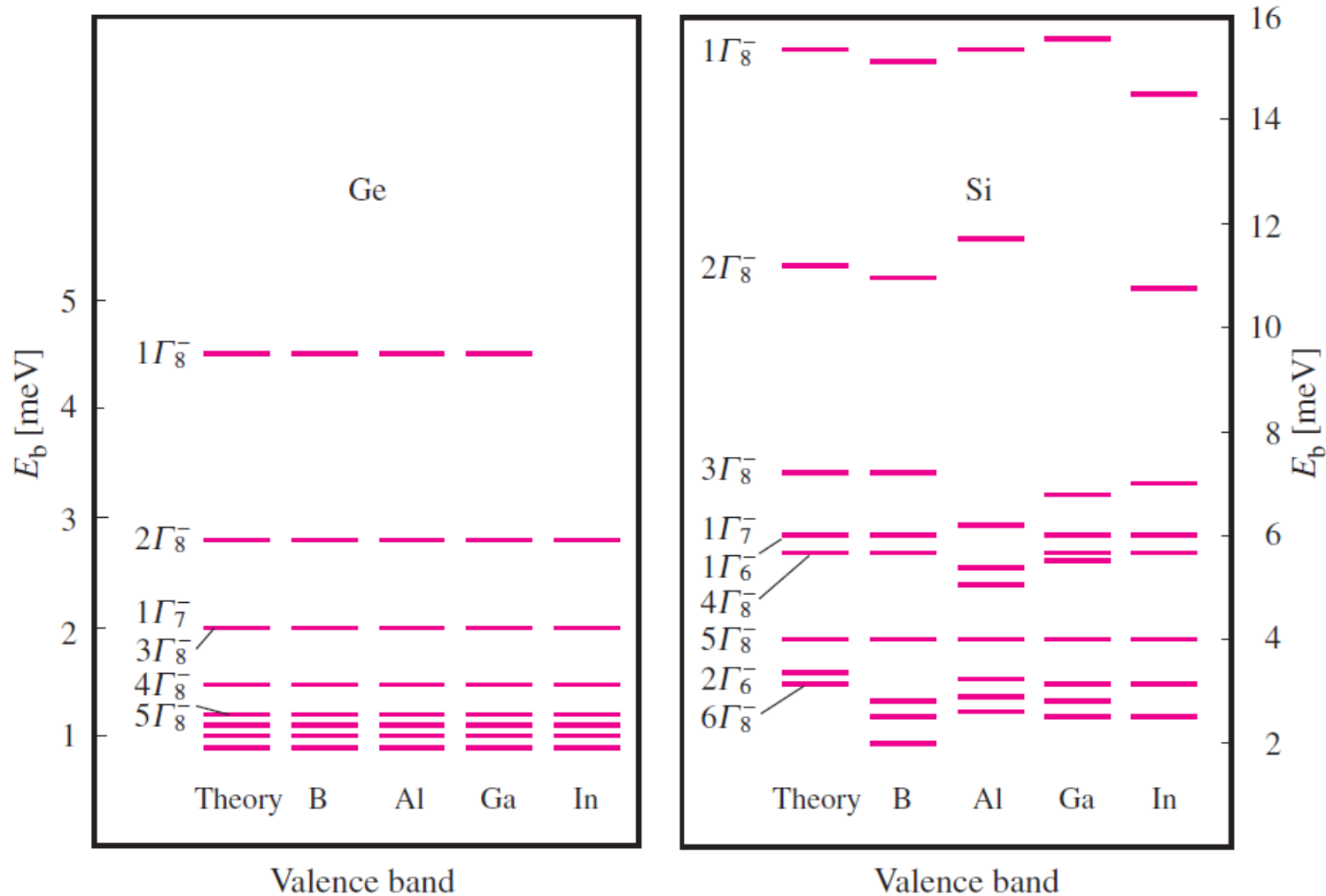
Works quite well for s-like conduction bands.

Complications due to p-like VB, anisotropic bands (Ge, Si, GaP)

Shallow (hydrogenic) defect: Donors in Si and Ge

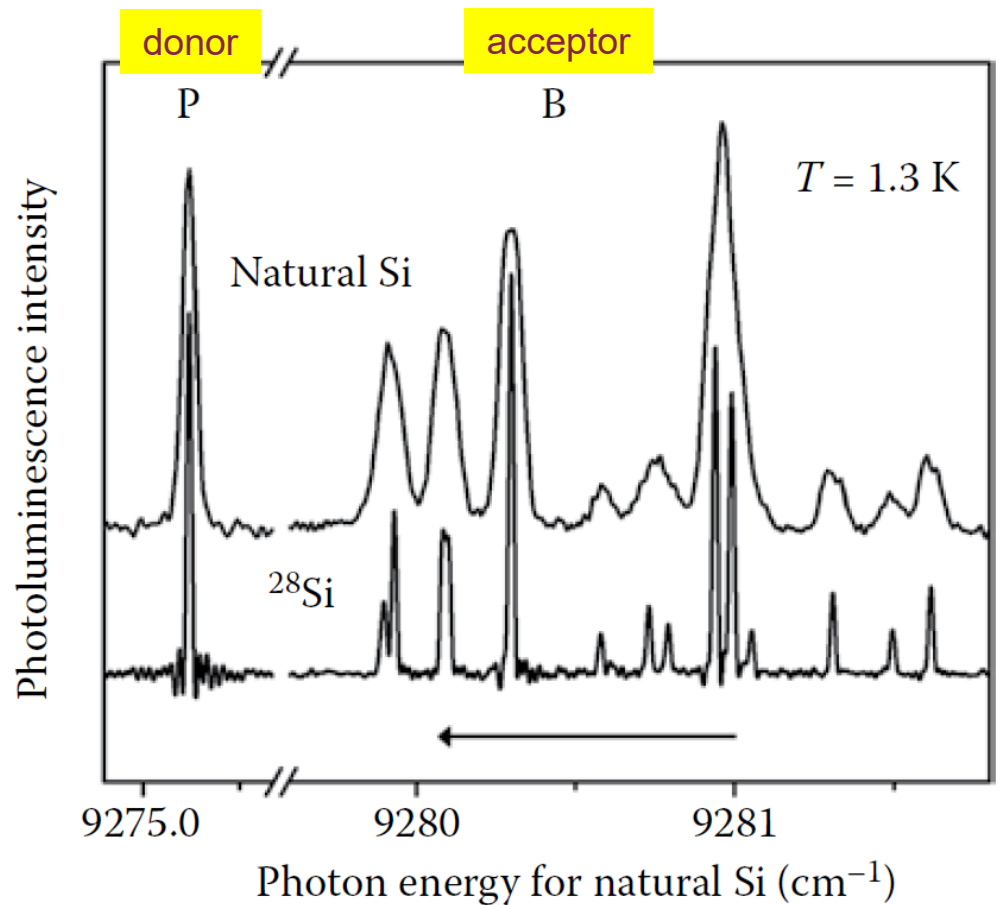
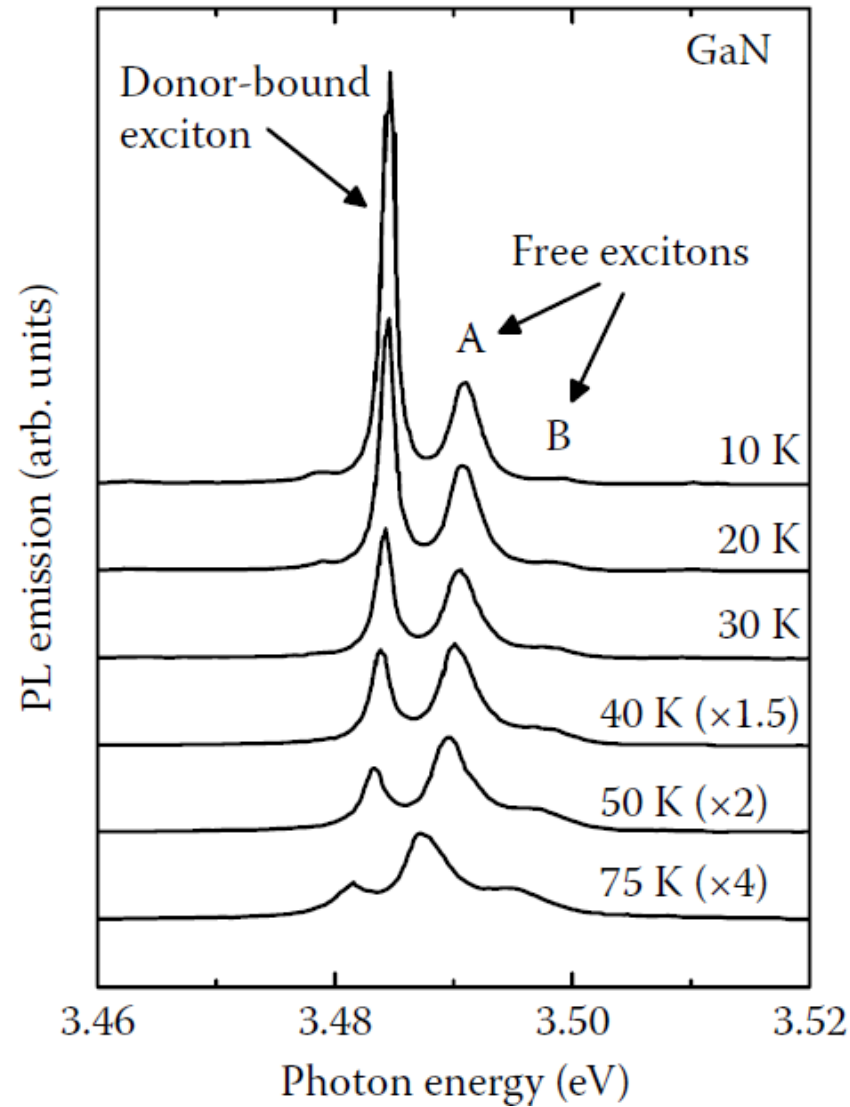


Shallow (hydrogenic) defect: Acceptors in Si and Ge



Need to include valence band warping (Luttinger theory) and other corrections.

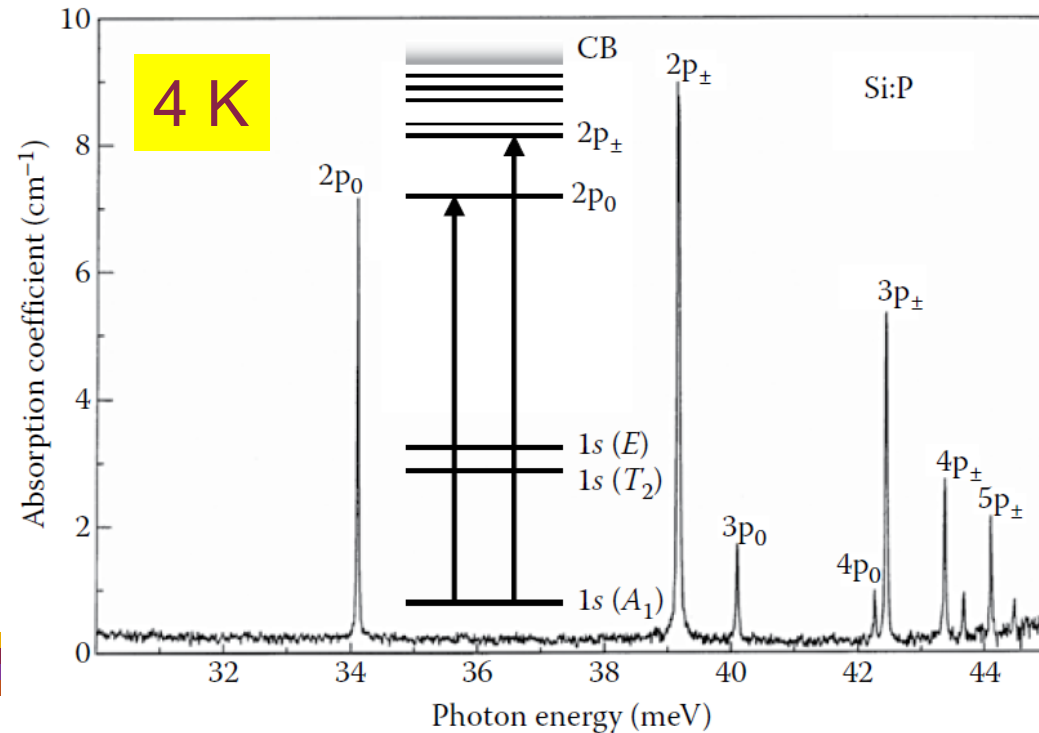
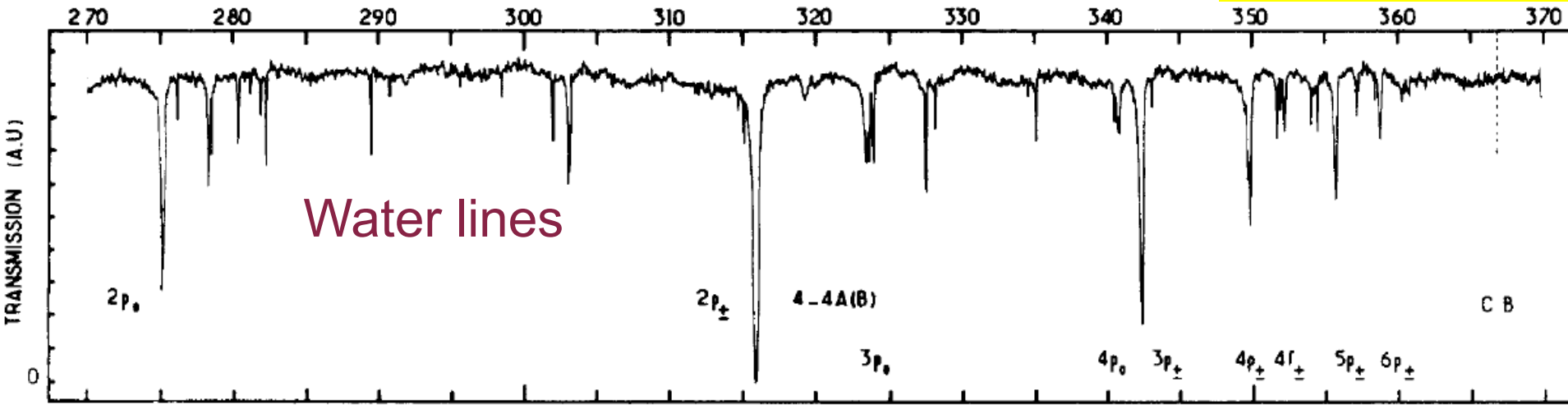
Free and bound excitons (photoluminescence)



Band to defect recombination shifted from free exciton energy.

Infrared defect spectroscopy: P donor in Si

Energy (cm⁻¹)

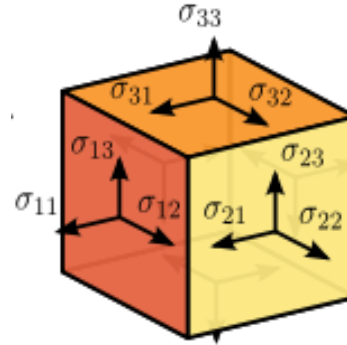


McClusky & Haller, Dopants and Defects
 A.K. Ramdas, PRB **23**, 2082 (1981)
 Pajot, Solid State Commun. **31**, 759 (1979)

Stress and Strain: Hooke's Law

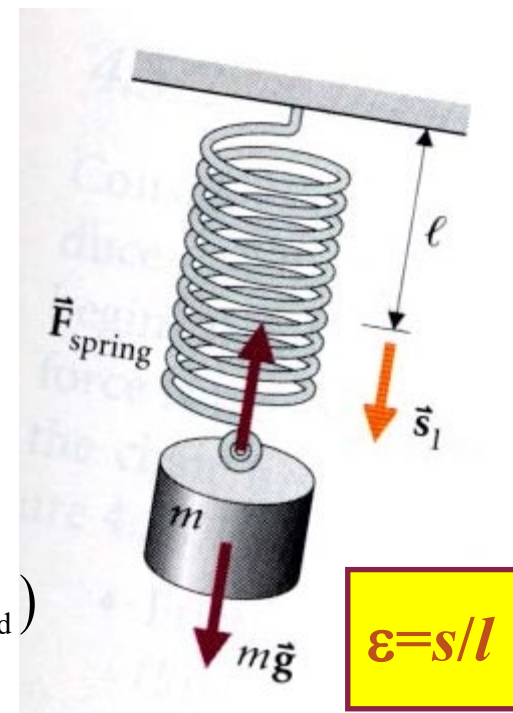
- **Stress:** Force per unit area (GPa), (1st rank tensor)

$$\mathbf{X} = \begin{pmatrix} X_{11} & X_{12} & X_{13} \\ X_{21} & X_{22} & X_{23} \\ X_{31} & X_{32} & X_{33} \end{pmatrix}$$



- **Strain:** Describes the response to the stress, (1st rank tensor)

$$\boldsymbol{\varepsilon} = \begin{pmatrix} \varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13} \\ \varepsilon_{21} & \varepsilon_{22} & \varepsilon_{23} \\ \varepsilon_{31} & \varepsilon_{32} & \varepsilon_{33} \end{pmatrix} = \vec{\nabla} \otimes (\vec{a}_{\text{strained}} - \vec{a}_{\text{unstrained}})$$



- Constituent relation (Hooke's law) between stress and strain: compliance tensor \mathbf{S} , stiffness tensor \mathbf{c} (6x6 2nd rank tensors)

$$\boldsymbol{\varepsilon} = \mathbf{S}\mathbf{X}$$

$$\mathbf{X} = \mathbf{c}\boldsymbol{\varepsilon}$$

Spring: Strain is $\boldsymbol{\varepsilon} = s/l$,
Relative change in length.

Landau & Lifshitz, *Elasticity Theory*
Yu & Cardona, *Fundamentals of Semiconductors*.



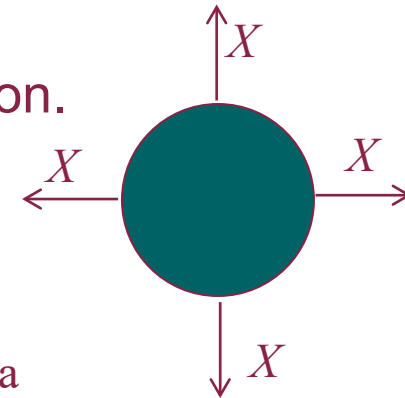
Biaxial Stress in Epitaxial Film on (001) substrate

- Stress:

Forces act along the wafer, but not along the growth direction.

- SiGe on Si: Compressive stress ($X < 0$)
- Si on SiGe: Tensile stress ($X > 0$)

$$\mathbf{X}_{\text{biaxial}} = \begin{pmatrix} X & 0 & 0 \\ 0 & X & 0 \\ 0 & 0 & 0 \end{pmatrix}$$



- Resulting Strain:

$$\boldsymbol{\epsilon}_{\text{biaxial}} = \begin{pmatrix} (S_{11} + S_{12})X & 0 & 0 \\ 0 & (S_{11} + S_{12})X & 0 \\ 0 & 0 & 2S_{12}X \end{pmatrix} = \begin{pmatrix} \epsilon_{\parallel} & 0 & 0 \\ 0 & \epsilon_{\parallel} & 0 \\ 0 & 0 & \epsilon_{\perp} \end{pmatrix}$$

$$\epsilon_{\parallel} = 0.553 X / 10^{11} \text{ Pa}$$

tensile in-plane strain

$$\epsilon_{\perp} = -0.979 X / 10^{11} \text{ Pa}$$

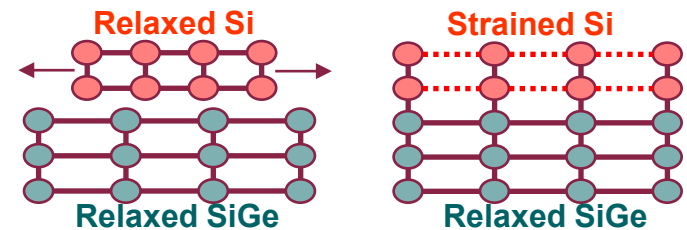
compressive vertical strain

- Hydrostatic strain component: tensile (>0), softens phonons, gap decreases

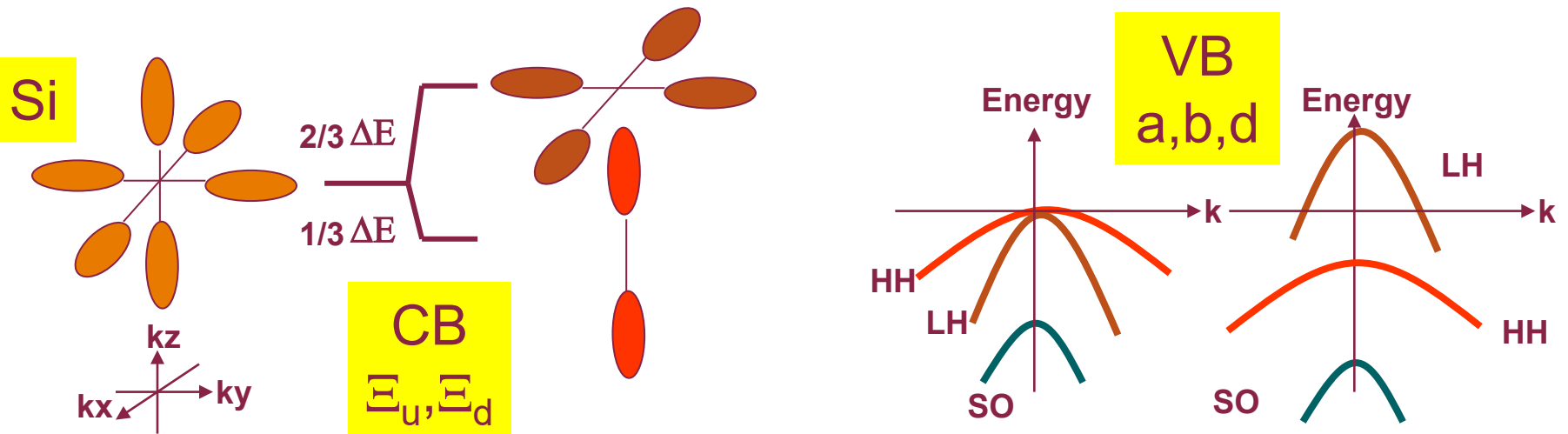
$$\epsilon_H = (2\epsilon_{\perp} + \epsilon_{\parallel}) / 3$$

- (100) Shear strain component: compressive (<0), Splits phonons and bands into singlet/doublet. Selection rules!

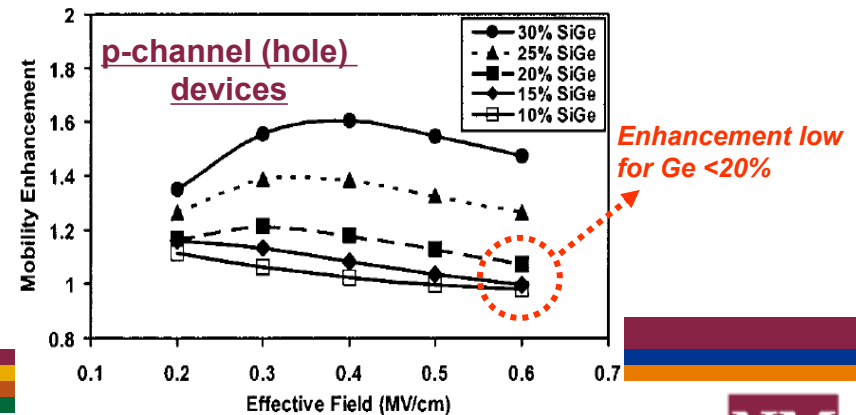
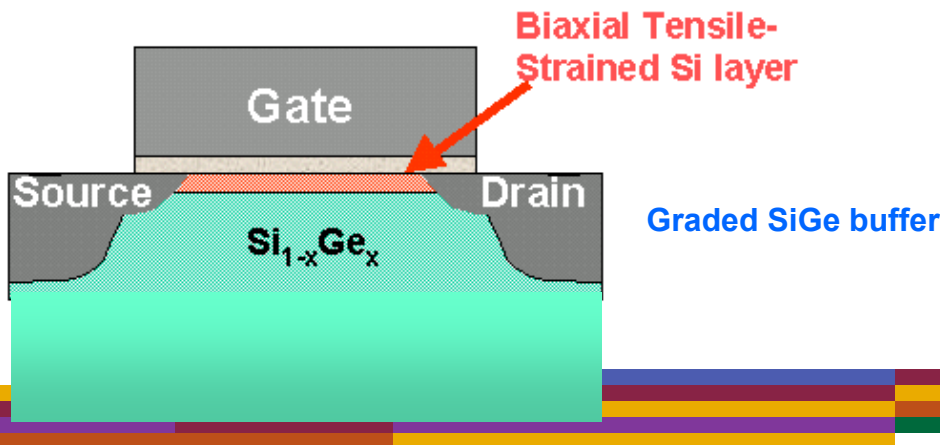
$$\epsilon_S = (\epsilon_{\parallel} - \epsilon_{\perp}) / 3$$



Thin Strained Silicon Layers for CMOS



- Biaxial tension lowers band gaps, reduces effective masses, and splits bands/valleys (reduced intervalley and inter-valence band scattering).
- Silicon under biaxial tension has higher electron and hole mobilities and therefore offers better transport properties than bulk silicon.



Deformation potentials

PHYSICAL REVIEW

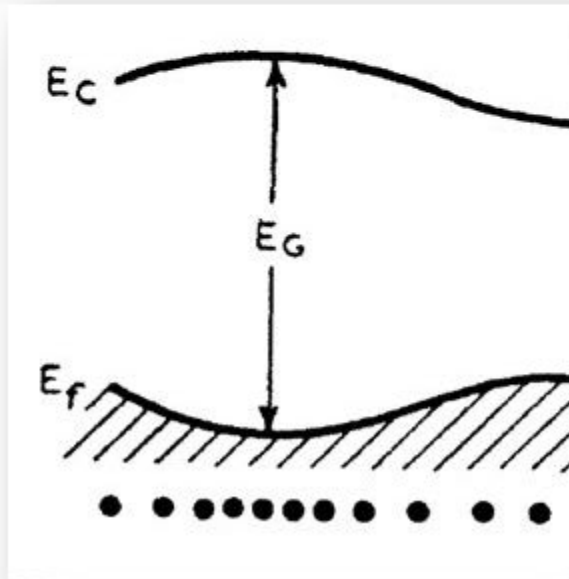
VOLUME 80, NUMBER 1

OCTOBER 1, 1950

Deformation Potentials and Mobilities in Non-Polar Crystals

J. BARDEEN AND W. SHOCKLEY
Bell Telephone Laboratories, Murray Hill, New Jersey
(Received May 31, 1950)

$$a = \frac{1}{V} \frac{\partial E_g}{\partial V}$$
$$\Delta E_g = a \text{Tr}(\epsilon)$$



$$a_V^{\Gamma-\Gamma} = \frac{\partial E_g^{\Gamma-\Gamma}}{\partial \ln V}$$

$$a_P^{\Gamma-\Gamma} = - \left(\frac{1}{B} \right) a_V^{\Gamma-\Gamma}$$

Summary

- Quantum confinement and Heisenberg uncertainty principle
- Growth of quantum structures
- **Carbon nanostructures, two-dimensional materials**
- **Electronic states, quantum well absorption and emission**
- **Intersubband transitions**
- **Metamaterials and metasurfaces**

- **Defects**
- **Transition metal and rare earth impurities in insulators**
- **Shallow defects in semiconductors**

- **Stress and strain, deformation potentials**